# Fluctuations of the first particle in exclusion processes 

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## Motivations

We consider continuous-time exclusion processes on $\mathbb{Z}$,

starting from the step initial condition


Under mild hypotheses, we expect that for $\kappa \in\left(0, \kappa^{*}\right)$,

$$
\frac{x_{\lfloor\kappa t\rfloor}-c t}{\sigma t^{1 / 3}} \Longrightarrow-\mathscr{L}_{\mathrm{GUE}}
$$

the Tracy-Widom GUE distribution.

## Question

Is the behaviour of $x_{1}(t)$ universal as well?

## Answer: NO

TASEP:


By the CLT, we have

$$
\frac{x_{1}(t)-t}{\sqrt{t}} \Longrightarrow \mathscr{N}
$$

The same limit theorem holds for any totally asymmetric exclusion processes.
ASEP: Let $R>L>0, R+L=1$ be asymmetry parameters


Theorem (Tracy-Widom 2009)

$$
\frac{x_{1}(t)-(R-L) t}{\sigma \sqrt{t}} \Longrightarrow \mathscr{X},
$$

where $\mathscr{X}$ is not a Gaussian. $\mathbb{P}(\mathscr{X} \leqslant x)=\operatorname{det}(I-K)_{\mathbb{L}^{2}(x, \infty)}$ where

$$
K(x, y)=\frac{R}{\sqrt{2 \pi}} e^{-\left(R^{2}+L^{2}\right) \frac{x^{2}+y^{2}}{4}+R L x y} .
$$

## MADM

The Multi-particle Asymmetric Diffusion Model (Sasamoto-Wadati 1998) is another exactly solvable partially asymmetric exclusion process.
Fix a parameter $q \in(0,1)$, asymmetry parameters $R>L>0, R+L=1$. The particle at $x_{n}(t)$ jumps to

- $x_{n}(t)+j$ at rate $\frac{R}{\left[j j_{q}-1\right.}$ for any $j \in\left\{1, \ldots, x_{n-1}(t)-x_{n}(t)-1\right\}$,
- $x_{n}(t)-j$ at rate $\frac{L}{[j]_{q}}$ for any $j \in\left\{1, \ldots, x_{n}(t)-x_{n+1}(t)-1\right\}$,
where the $q$ - deformed integer $[j]_{q}$ is given by

$$
\begin{gathered}
{[j]_{q}=1+q+\cdots+q^{j-1},} \\
{[j]_{q^{-1}}=1+q^{-1}+\cdots+q^{-j+1} .}
\end{gathered}
$$



## Limit Theorem

## Theorem (B.-Corwin 2014)

There exist constants $c, \sigma, L^{*}$ such that for $0<L<L^{*}$

$$
\frac{x_{1}(t)-c t}{\sigma t^{1 / 3}} \Longrightarrow-\mathscr{L}_{\mathrm{GUE}}
$$

The result should hold with $L^{*}=1 / 2$. The first particle behaves as in the bulk. Indeed, one can prove the one-point asymptotics predicted by KPZ universality,

## Theorem (B.-Corwin 2014)

There exist constants $c(\kappa), \sigma(\kappa), L^{*}, \kappa^{*}$ such that for $0 \leqslant L<L^{*}$ and $\kappa \in\left(0, \kappa^{*}\right)$,

$$
\frac{x_{\lfloor\kappa t\rfloor}(t)-c(\kappa) t}{\sigma(\kappa) t^{1 / 3}} \Longrightarrow-\mathscr{L}_{\mathrm{GUE}}
$$

## Why so different than ASEP?

Let $\rho(x):=$ density of particles around $x t$ at time $t$ as $t$ goes to infinity.
TASEP


ASEP



## Universality?

## Question

For exclusion processes such that the density around the first particle is positive, are the $t^{1 / 3}$ scaling and GUE distribution universal?

In order to test the universality, one needs at least one other such process.

## Question

When is the density of particles positive around the first particle?
The density profile has a jump discontinuity when the drift (average speed of a tagged particle) is not decreasing as a function of the local density.

## Hydrodynamic limit

- Assume that there exists a family of translation invariant stationary measures indexed by the average density of particles $\varrho$.
- Define the flux $j(\varrho)$ as the expected (for that measure) number of particles crossing a given bound per unit of time, counted algebraically.
- Assume that the following limit exists

$$
\rho(x, t):=\lim _{\tau \rightarrow \infty} \mathbb{P}(\text { There is a particle at site } x \tau \text { at time } t \tau) .
$$

It satisfies the conservation equation

$$
\frac{\partial}{\partial t} \rho(x, t)+\frac{\partial}{\partial x} j(\rho(x, t))=0
$$

heuristic result: Let $\varrho_{0}$ be the density of particles around the first particle. The density profile is discontinuous at the first particle (i.e. $\left.\varrho_{0}>0\right)$ when the function $j(\varrho) / \varrho$ is not decreasing. Actually $\varrho_{0}$ locally maximizes the drift, $j(\varrho) / \varrho$.

## Heuristic proof

Assume $\varrho_{0}>0$.
(1) On the one hand, the macroscopic position of the first particle is its drift $j\left(\rho_{0}\right) / \varrho_{0}$.
(2) On the other hand the characteristics method (applied to the conservation PDE) yields a function $\pi(\varrho)$ s.t.

$$
\begin{equation*}
\rho(\pi(\varrho) t, t)=\varrho . \tag{1}
\end{equation*}
$$

i.e. $\pi(\rho)$ is the macroscopic position where particles have a local density $\varrho$. Differentiating (1) yields

$$
\pi(\varrho)=\frac{\partial j(\varrho)}{\partial \varrho}=j^{\prime}(\varrho)
$$

Combining (1) and (2), we have that

$$
j^{\prime}\left(\varrho_{0}\right)=\left.\frac{j\left(\varrho_{0}\right)}{\varrho_{0}} \Longrightarrow \frac{\mathrm{~d}}{\mathrm{~d} \varrho} \frac{j(\varrho)}{\varrho}\right|_{\varrho=\varrho_{0}}=0
$$

which suggests that $\varrho_{0}$ is a maximizer of $j(\varrho) / \varrho$.

## Facilitated TASEP

## Question

For exclusion processes such that the density around the first particle is positive, are the $t^{1 / 3}$ scaling and GUE distribution universal?

We consider the Facilitated TASEP (FTASEP): the particle at $x_{n}(t)$ moves by +1 at rate 1 provided that

- the site $x_{n}(t)+1$ is empty (exclusion),
- the site $x_{n}(t)-1$ is occupied (facilitation).


Introduced in physics literature, Basu-Mohanty 2009, and studied further by Gabel-Krapivsky-Redner 2010. The flux

$$
j(\rho)=\frac{(1-\rho)(2 \rho-1)}{\rho} \mathbb{1}_{\rho>1 / 2}
$$

is such that $j(\rho) / \rho$ has a maximum for $\rho=2 / 3$.

The density profile is given by

$$
\rho(x)=\frac{1}{\sqrt{2+x}} \text { for } x \in(-1,1 / 4) .
$$



## Theorem (Baik-B.-Corwin-Suidan)

$$
\frac{x_{1}(t)-t / 4}{2^{-4 / 3} t^{1 / 3}} \Longrightarrow-\mathscr{L}_{\mathrm{GSE}}
$$

where $\mathscr{L}_{\text {GSE }}$ is the Tracy-Widom GSE distribution.
The FTASEP is in the KPZ universality class in the sense that

## Theorem (Baik-B.-Corwin-Suidan)

For all $r \in(0,1)$, there exist (explicit) constants $\pi(r), \sigma(r)$ such that

$$
\frac{x_{\lfloor r t\rfloor}(t)-t \pi(r)}{\sigma(r) t^{1 / 3}} \Longrightarrow-\mathscr{L}_{\mathrm{GUE}}
$$

as the KPZ scaling theory predicts.

## Proofs

- MADM: it can be studied via a method initially designed by Borodin-Corwin-Sasamoto 2012 for the $q$-TASEP and ASEP, using Markov duality and Bethe ansatz.
- FTASEP: the solvability comes from a coupling with last passage percolation on a half-quadrant.


## FTASEP and OpenTASEP



We use first a coupling between the FTASEP and a TASEP on a semi-infinite lattice with a source at the origin (we call it the OpenTASEP)


Define the current at site $x$ by

$$
N_{x}(t)=\#\{i \geqslant x \mid \text { site } i \text { is occupied }\} .
$$

## The coupling

Consider the gaps between consecutive particles in the FTASEP

$$
g_{i}(t):=x_{i}(t)-x_{i+1}(t)-1
$$

For all $i \geqslant 1$, the rules of the dynamics implies that $g_{i} \in\{0,1\}$.


The current at site $n$ in the OpenTASEP corresponds to the number of jumps done by the $n$th particle in FTASEP, i.e. $x_{n}(t)+n$.

## Proposition

We have the equality in law of the processes

$$
\left\{x_{n}(t)+n\right\}_{n \geqslant 1, t \geqslant 0}=\left\{N_{n}(t)\right\}_{n \geqslant 1, t \geqslant 0} .
$$

Let us see how this works dynamically


- gray: particle that cannot move,
- black: particle that can move.

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## Last passage percolation

- Let $w_{i j}$ a family of i.i.d. exponential random variables.
- Consider up-right paths $\pi$ from the box $(1,1)$ to ( $n, m$ ) in the half quadrant. We define the last passage percolation time $H(n, m)$ by

$$
H(n, m)=\max _{\pi} \sum_{(i, j) \in \pi} w_{i j} .
$$

## Lemma

If $w_{i j} \sim \operatorname{Exp}(1)$,

$$
\mathbb{P}\left(N_{n}(t) \leqslant x\right)=\mathbb{P}(H(n+x-1, x) \geqslant t)
$$


$x_{1}(t)$ in FTASEP corresponds to $H(n, n)$.

## Passage-times on the diagonal

LPP in a half-quadrant has first been studied by Baik and Rains (2001) with Geometric weights. In the model with exponential weights, we find similar limit theorems.

## Theorem (Baik-B.-Corwin-Suidan)

Assume that $w_{i j} \sim \operatorname{Exp}(1)$ for $i>j$ and $w_{i i} \sim \operatorname{Exp}(\alpha)$ for some parameter $\alpha>0$.

- When $\alpha>1 / 2$,

$$
\frac{H(n, n)-4 n}{2^{4 / 3} n^{1 / 3}} \Longrightarrow \mathscr{L}_{\mathrm{GSE}}
$$

(implies the GSE limit theorem for $x_{1}(t)$ in FTASEP, corresponding to $\alpha=1$.)

- When $\alpha=1 / 2$,

$$
\frac{H(n, n)-4 n}{2^{4 / 3} n^{1 / 3}} \Longrightarrow \mathscr{L}_{\mathrm{GOE}}
$$

- When $\alpha<1 / 2$,

$$
\frac{H(n, n)-c n}{c^{\prime} n^{1 / 2}} \Longrightarrow \mathscr{N}
$$

The parameter $\alpha$ corresponds to the rate of the first particle in the FTASEP.

## Away from the diagonal: KPZ typical behaviour

The fluctuations away from the diagonal have first been studied by Sasamoto-Imamura 2004 - for the discrete PNG model. In the model with exponential weights, we have

## Theorem (Baik-B.-Corwin-Suidan)

For $\kappa \in(0,1)$ and $\alpha>\sqrt{\kappa} /(1+\sqrt{\kappa})$,

$$
\frac{H(n, \kappa n)-(1+\sqrt{\kappa})^{2} n}{\sigma n^{1 / 3}} \Longrightarrow \mathscr{L}_{\mathrm{GUE}}
$$

(implies the GUE limit theorem for $x_{(1-\kappa) t}$ in FTASEP)

## Proofs?

(I) LPP in a half-quadrant is a marginal of a Pfaffian Schur process.
(II) By a theorem of Borodin-Rains 2005, it is hence a Pfaffian point process, with explicit correlation kernel.
(III) Saddle-point analysis of the correlation kernel yields the various limit theorems (in progress).

For integer partitions $\lambda_{1} \geqslant \lambda_{2} \geqslant \ldots$, and $\mu_{1} \geqslant \mu_{2} \geqslant \ldots$, we will consider skew-Schur functions

$$
s_{\lambda / \mu}=\operatorname{det}\left[h_{\lambda_{i}-\mu_{j}+j-i}\right]_{i, j},
$$

where $h_{k}$ are complete homogeneous symmetric functions

$$
h_{k}(x)=\sum_{i_{1} \leqslant \cdots \leqslant i_{k}} x_{i_{1}} \ldots x_{i_{k}} .
$$

We also define

$$
\tau_{\lambda}=\sum_{\kappa^{\prime} \text { even }} s_{\lambda / \kappa}=\operatorname{Pf}[\ldots]
$$

where $\kappa^{\prime}$ even means that $\kappa_{1}=\kappa_{2} \geqslant \kappa_{3}=\kappa_{4} \geqslant \ldots$.

Schur process


Consider a path $\gamma$ as on the left
$\rightarrow$ vertex $v \mapsto \lambda^{v}$ a random partition,

- edge $e \mapsto \rho_{e}$ a set of variables. (More generally a specialization of the symmetric functions).

The Schur process (Okounkov-Reshetikhin 2003) is a probability measure on the sequence of partitions $\lambda:=\left(\lambda^{v}\right)_{v \in \gamma}$ such that

$$
\mathbb{P}(\lambda)=\frac{1}{Z} \prod_{e \in \gamma} \text { weight }(e)=\frac{1}{Z} \operatorname{det}[\ldots],
$$

where

$$
\text { weight }\left(e={v^{\prime}}^{\leftarrow} v\right)=s_{\lambda^{v} / \lambda^{v^{\prime}}}\left(\rho_{e}\right) \text { and weight }\left(e=\uparrow_{v}^{v^{\prime}}\right)=s_{\lambda^{v^{\prime} / \lambda v^{v}}}\left(\rho_{e}\right)
$$

## Pfaffian Schur process



Consider a path $\gamma$ as on the left

- vertex $v \mapsto \lambda^{v}$ a random partition,
- edge $e \mapsto \rho_{e}$ a set of variables.
- Denote $\rho_{\circ}$ and $\lambda_{\circ}$ the specialization and the partition on the diagonal.

The Pfaffian Schur process is a probability measure on the sequence of partitions $\lambda:=\left(\lambda^{v}\right)_{v \in \gamma}$ such that

$$
\left.\mathbb{P}(\lambda)=\frac{1}{Z} \tau_{\lambda_{0}}\left(\rho_{\circ}\right) \prod_{e \in \gamma} \text { weight }(e)=\frac{1}{Z} \operatorname{Pf[} \ldots\right],
$$

where the weight of off-diagonal edges are chosen as in the Schur process.

## Geometric last passage percolation

Assume that all $\rho_{e}=\{\sqrt{q}\}$, and $\rho_{\circ}=\{c\}$. Then for $0<n_{1} \leqslant \cdots \leqslant n_{k}$, $m_{1} \geqslant \cdots \geqslant m_{k}$, with $n_{i} \geqslant m_{i}$,

$$
\left(\lambda_{1}^{\left(n_{1}, m_{1}\right)}, \ldots, \lambda_{1}^{\left(n_{k}, m_{k}\right)}\right) \stackrel{(d)}{=}\left(G\left(n_{1}, m_{1}\right), \ldots, G\left(n_{k}, m_{k}\right)\right)
$$

where the family of random variables $G(n, m)$ satisfies the recursion

$$
\left\{\begin{array}{l}
G(n, m)=\max \{G(n-1, m), G(n, m-1)\}+\operatorname{Geom}(q) \text { for } n>m \\
G(n, n)=G(n, n-1)+\operatorname{Geom}(c \sqrt{q}) .
\end{array}\right.
$$

As the geometric distribution converges to the exponential,

## Proposition

If we set $c=\sqrt{q}(1+(\alpha-1)(q-1))$, then as $q \rightarrow 1$,

$$
\left\{(1-q) G\left(n_{i}, m_{i}\right)\right\}_{i=1}^{k} \Longrightarrow\left\{H\left(n_{i}, m_{i}\right)\right\}_{i=1}^{k}
$$

where $H(n, m)$ are the passage times in LPP with exponential weights on a half quadrant (and parameter $\alpha$ on the diagonal).

## Pfaffian Point process

A random configuration $X \subset \mathbb{X}$ (state space) is a Pfaffian point process if one can write the correlation function as

$$
\rho(Y)=\mathbb{P}(Y \subset X)=\operatorname{Pf}[K(x, y)]_{x, y \in Y},
$$

where

$$
K(x, y)=\left(\begin{array}{ll}
K_{11}(x, y) & K_{12}(x, y) \\
K_{21}(x, y) & K_{22}(x, y)
\end{array}\right)
$$

is a skew-symmetric matrix indexed by elements in $\mathbb{X}$; called the correlation kernel.
The gap probabilities are given by Fredholm Pfaffians

$$
\mathbb{P}(\text { no point in } Y)=\operatorname{Pf}(J-K)_{\mathbb{L}^{2}(Y)}
$$

where

$$
\operatorname{Pf}(J-K)_{\mathbb{L}^{2}(Y)}:=1+\sum_{k=1}^{\infty} \frac{(-1)^{k}}{k!} \int_{Y} \mathrm{~d} x_{1} \ldots \int_{Y} \mathrm{~d} x_{k} \operatorname{Pf}\left[K\left(x_{i}, x_{j}\right)\right]_{i, j=1}^{k}
$$

## The Pfaffian Schur process is Pfaffian

## Theorem (Borodin-Rains 2005)

For $0<n_{1} \leqslant \cdots \leqslant n_{k}, m_{1} \geqslant \cdots \geqslant m_{k}$, with $n_{i} \geqslant m_{i}$, the Pfaffian Schur process is Pfaffian in the sense that

$$
\left(1, \lambda_{i}^{\left(n_{1}, m_{1}\right)}-i\right)_{i \geqslant 1} \cup \cdots \cup\left(k, \lambda_{i}^{\left(n_{k}, m_{k}\right)}-i\right)_{i \geqslant 1} \in \mathbb{X}=\{1, \ldots, k\} \times \mathbb{Z}
$$

is a Pfaffian point process with an explicit correlation kernel $K$.
The variables $G\left(n_{i}, m_{i}\right) \stackrel{(d)}{=} \lambda_{1}^{\left(n_{i}, m_{i}\right)}$ are extremal points in the Pfaffian point process, so that

$$
\mathbb{P}\left(G\left(n_{1}, m_{1}\right) \leqslant h_{1}, \ldots, G\left(n_{k}, m_{k}\right) \leqslant h_{k}\right)=\operatorname{Pf}(J-K)_{\mathbb{L}^{2}(\ldots)} .
$$

Finally, sending $q \rightarrow 1$ yields the probability distribution of passage times in exponential LPP on the half-quadrant.

In the limit, the state space becomes $\{1, \ldots, k\} \times \mathbb{R}$.

## Proposition (Baik-B.-Corwin-Suidan)

For $0<n_{1}<\cdots<n_{k}, m_{1}>\cdots>m_{k}$ with $n_{i}>m_{i}, h_{1}, \ldots, h_{k}>0$

$$
\mathbb{P}\left(H\left(n_{1}, m_{1}\right) \leqslant h_{1}, \ldots, H\left(n_{k}, m_{k}\right) \leqslant h_{k}\right)=\operatorname{Pf}\left(J-K^{\exp }\right)_{\mathbb{L}^{2}}\left(\Delta_{k}\left(h_{1}, \ldots, h_{k}\right)\right) .
$$

where

$$
\Delta_{k}\left(h_{1}, \ldots, h_{k}\right)=\left\{(i, x) \in \mathbb{Z} \times \mathbb{R} \mid x>h_{i}\right\},
$$

and the kernel $K$ is given by

$$
\begin{aligned}
& K_{11}^{\exp }(i, x ; j, y)=\frac{1}{(2 \mathbf{i} \pi)^{2}} \int_{\infty \infty^{- \text {- } \pi / 3}}^{\infty e^{\mathrm{i} \pi / 3}} \mathrm{~d} z \int_{\infty e^{-\mathrm{i} \pi / 3}}^{\infty} \mathrm{d} w \frac{z-w}{\infty \mathrm{i} \pi / 3} e^{-x z-y w} \\
& \frac{(1+2 z)^{n_{i}}(1+2 w)^{n_{j}}}{(1-2 z)^{m_{i}}(1-2 w)^{m_{j}}}(2 z+2 \alpha-1)(2 w+2 \alpha-1),
\end{aligned}
$$

where the contours pass to the right of 0 , and we have formulas of a similar taste for $K_{12}$ and $K_{22}$.

Since the GSE/GOE/GUE distribution functions can be written as a Fredholm Pfaffian, one concludes by asymptotic analysis of the above formula.

## Summary

We have seen that

- The fluctuations of the first particle in exclusion processes are not universal.
- For the FTASEP, we find the GSE Tracy-Widom distribution.
- This is proved via a coupling with Last Passage Percolation in a half-quadrant.
- Which can be studied exhaustively via Pfaffian Schur Processes, when the weights are geometric or exponential.


## Outlook

## Further directions

- One can play with parameters in LPP, proving phase transitions and studying crossover distributions.
- There are other marginals of the Pfaffian Schur process (other particle dynamics, symmetric plane partitions...).
- Pfaffian Schur processes can be leveraged to Pfaffian Macdonald processes, leading to positive temperature models.


## Questions

- In presence of a jump discontinuity, can one prove the $t^{1 / 3}$ behaviour in general?
- Can one understand the geometric behaviour of the geodesic in LPP ? give a probabilistic interpretation of the phase transition? Compare to the slow bond problem.


## Thank you

## Proofs for MADM

## MADM



The limit theorem follows from

- A Markov duality between the MADM exclusion process and a zero range analogue, so that for $\vec{n}=n_{1} \geqslant n_{2} \geqslant \cdots \geqslant n_{k}$, the function

$$
(t, \vec{n}) \mapsto \mathbb{E}\left[\prod_{i=1}^{k} q^{x_{n_{i}}(t)}\right]
$$

satisfies a closed system of differential equations (Kolmogorov equation for the dual system).

- This system of ODEs is solvable via Bethe ansatz. It leads to contour integral formulas for the moments of $q^{x_{n}(t)}$.

$$
\begin{aligned}
\mathbb{E}\left[\prod_{i=1}^{k} q^{x_{n_{i}}(t)+n_{i}}\right] & =\frac{(-1)^{k} q^{\frac{k(k-1)}{2}}}{(2 \pi i)^{k}} \oint_{\gamma_{1}} \cdots \oint_{\gamma_{k}} \prod_{1 \leq A<B \leq k} \frac{z_{A}-z_{B}}{z_{A}-q z_{B}} \\
& \times \prod_{j=1}^{k}\left(\frac{1-q z_{j}}{1-z_{j}}\right)^{n_{j}} \exp \left((q-1) t\left(\frac{R z_{j}}{1-q z_{j}}-\frac{L z_{j}}{1-z_{j}}\right)\right) \frac{d z_{j}}{z_{j}\left(1-q z_{j}\right)},
\end{aligned}
$$

where the integration contours $\gamma_{1}, \ldots, \gamma_{k}$ are nested in order to enclose all poles except 0 and $1 / q$.

- The moments do characterize the distribution of $x_{n}(t)$. One can take the moment generating function and form the ( $q$-deformed) Laplace transform of $q^{x_{n}(t)}$.
- Rearranging terms as in a Fredholm determinant expansion, a saddle-point asymptotic analysis yields the GUE limit theorem.


## Dynamics on the Pfaffian Schur Process

We define dynamics preserving the Pfaffian Schur processes that correspond to LPP in a half quadrant. We make a path $\gamma$ grow as follows




At each stage we consider a Pfaffian Schur process indexed by the path. We update the partitions where the path has changed according to Markov transition kernels.

When the path grows by one box from a corner formed by partitions $\kappa, \mu$ and $v$, we update according to some transition kernel

where we need that

$$
\sum_{\mu} s_{\kappa / \mu}\left(\rho_{2}\right) s_{v / \mu}\left(\rho_{1}\right) \mathscr{U}_{\rho_{1}, \rho_{2}}^{\llcorner }(\pi \mid v, \mu, \kappa)=\text { const. } s_{\pi / \kappa}\left(\rho_{1}\right) s_{\pi / v}\left(\rho_{2}\right)
$$

so that the Pfaffian Schur structure is preserved. const is a normalization constant depending only on the specializations $\rho_{1}, \rho_{2}$. We choose

$$
\mathscr{U}_{\rho_{1}, \rho_{2}}^{\llcorner }(\pi \mid v, \mu, \kappa)=\mathscr{U}_{\rho_{1}, \rho_{2}}^{\llcorner }(\pi \mid v, \kappa)=\frac{s_{\pi / v}\left(\rho_{2}\right) s_{\pi / \kappa}\left(\rho_{1}\right)}{\sum_{\lambda} s_{\lambda / v}\left(\rho_{2}\right) s_{\lambda / \kappa}\left(\rho_{1}\right)} .
$$

This corresponds to so-called "push-block" dynamics in the usual (determinantal) Schur process.

Similarly, when the path grows by a half-box along the diagonal, we update according to

where we need that

$$
\sum_{\mu} s_{\kappa / \mu}\left(\rho_{1}\right) \tau_{\mu}\left(\rho_{\circ}\right) \mathscr{U}_{\rho_{\circ}, \rho_{1}}^{\llcorner }(\pi \mid \kappa, \mu)=\text { const. } s_{\pi / \kappa}\left(\rho_{1}\right) \tau_{\pi}\left(\rho_{\circ}\right)
$$

so that the Pfaffian Schur structure is preserved.
We choose

$$
\mathscr{U}_{\rho_{\circ}, \rho_{1}}^{\angle}(\pi \mid \kappa, \mu)=\mathscr{U}_{\rho_{\circ}, \rho_{1}}^{\angle}(\pi \mid \kappa)=\operatorname{const} \frac{\tau_{\pi}\left(\rho_{\circ}\right) s_{\pi / \kappa}\left(\rho_{1}\right)}{\tau_{\kappa}\left(\rho_{\circ}, \rho_{1}\right)} .
$$

## First coordinate marginal

Assume that all $\rho_{e}$ are specializations into a single variable $\rho_{e}=\sqrt{q}$, and $\rho_{\circ}=c$. Then we have that

$$
s_{\lambda / \mu}\left(\rho_{e}\right)=\mathbb{1}_{\mu<\lambda}(\sqrt{q})^{\sum \lambda_{i}-\sum \mu_{i}}
$$

where

$$
\mu<\lambda \Longleftrightarrow \lambda_{1} \geqslant \mu_{1} \geqslant \lambda_{2} \geqslant \geqslant \mu_{2} \geqslant \ldots,
$$

and

$$
\tau_{\lambda}\left(\rho_{\circ}\right)=c^{\lambda_{1}-\lambda_{2}+\lambda_{3}-\lambda_{4}+\ldots} .
$$

- Under the transition operator $\mathscr{U}^{\llcorner }(\pi \mid v, \kappa)$,

$$
\pi_{1}=\max \left\{v_{1}, \kappa_{1}\right\}+\operatorname{Geom}(q) .
$$

- Under the transition operator $\mathscr{U}^{\perp}(\pi \mid \kappa)$,

$$
\pi_{1}=\kappa_{1}+\operatorname{Geom}(q)
$$

## Geometric last passage percolation

It implies that for $0<n_{1} \leqslant \cdots \leqslant n_{k}, m_{1} \geqslant \cdots \geqslant m_{k}$, with $n_{i} \geqslant m_{i}$,

$$
\left.\left(\lambda_{1}^{\left(n_{1}, m_{1}\right)}, \ldots, \lambda_{1}^{\left(n_{k}, m_{k}\right)}\right)^{(d)} \stackrel{(G)}{=}\left(n_{1}, m_{1}\right), \ldots, G\left(n_{k}, m_{k}\right)\right)
$$

where the family of random variables $G(n, m)$ satisfies the recursion

$$
\left\{\begin{array}{l}
G(n, m)=\max \{G(n-1, m), G(n, m-1)\}+\operatorname{Geom}(q) \text { for } n>m \\
G(n, n)=G(n, n-1)+\operatorname{Geom}(q)
\end{array}\right.
$$

As the geometric distribution converges to the exponential,

## Proposition

If we set $\rho_{\circ}=c=\sqrt{q}(1+(\alpha-1)(q-1))$, then as $q \rightarrow 1$,

$$
\left\{(1-q) G\left(n_{i}, m_{i}\right)\right\}_{i=1}^{k} \Longrightarrow\left\{H\left(n_{i}, m_{i}\right)\right\}_{i=1}^{k}
$$

where $H(n, m)$ are the passage times in LPP with exponential weights on a half quadrant (and parameter $\alpha$ on the diagonal).

