

Riemann's second proof of the analytic continuation of the Riemann Zeta function

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Seminar on Modular Forms, Winter term 2006

1 Abstract

The Riemann zeta-function $\zeta(s)$ is defined by

$$\zeta(s) := \sum_{n=1}^{\infty} \frac{1}{n^s} \quad (1)$$

for $\operatorname{Re} s > 1$, but it is well known that there exists an analytic continuation onto the whole s -plane with a simple pole at $s = 1$. We want to derive this result using tools like the theta function

$$\theta(t) := \sum_{n=-\infty}^{\infty} e^{-\pi n^2 t}, \quad t > 0, \quad (2)$$

its Mellin transform

$$\int_0^{\infty} \theta(t) t^s \frac{dt}{t}, \quad s \in \mathbb{C},$$

the Gamma function

$$\Gamma(s) := \int_0^{\infty} e^{-t} t^s \frac{dt}{t}, \quad \operatorname{Re} s > 0, \quad (3)$$

and Fourier transformation.

This proof of the analytic continuation is known as the second Riemannian proof.

2 Some tools

2.1 The Gamma function

Remark: The Gamma function has a large variety of properties. Most of those we use are very well known, but we will provide all the proofs anyways.

Proposition 1: $\Gamma(s)$ satisfies the functional equation

$$\Gamma(s+1) = s\Gamma(s) \quad (4)$$

Proof:

$$\Gamma(s+1) = \int_0^\infty e^{-t} t^{s+1} \frac{dt}{t} = -e^{-t} t^s \Big|_0^\infty - \int_0^\infty -e^{-t} s t^s \frac{dt}{t} = s \Gamma(s)$$

□

Corollary: $\Gamma(s)$ has an analytic continuation on \mathbb{C} with simple poles at $s = 0, -1, -2, \dots$

Proof: Using (4) we can find values for $-1 < \operatorname{Re} s \leq 0$, except for $s = 0$. These values give us an analytic continuation for $\operatorname{Re} s > -1$ with a simple pole at $s = 0$. Of course we can repeat this step infinitely many times, with poles showing up at $s = 0, -1, -2, \dots$ □

Proposition 2: $\Gamma(s)$ satisfies the functional equation

$$\Gamma(s) \Gamma(1-s) = \frac{\pi}{\sin \pi s}$$

Proof: $\Gamma(s)$ has an other representation which equals (3):

$$\Gamma(s) = \frac{1}{s} \prod_{n=1}^{\infty} \frac{(1 + \frac{1}{n})^s}{1 + \frac{s}{n}}$$

With this representation and the Euler sine product we get

$$\begin{aligned} \Gamma(s) \Gamma(1-s) &= -s \Gamma(s) \Gamma(-s) \\ &= \frac{-s}{-s \cdot s} \prod_{n=1}^{\infty} \frac{1}{1 - \frac{s^2}{n^2}} \\ &= \frac{1}{s} \cdot \frac{\pi s}{\sin \pi s} \\ &= \frac{\pi}{\sin \pi s} \end{aligned}$$

□

Corollary: $\frac{1}{\Gamma(s)} = \frac{\sin \pi s}{\pi} \Gamma(1-s)$ is entire.

Proof: $\Gamma(1-s)$ has simple poles at $s = 1, 2, \dots$, but there $\sin \pi s$ has simple zeros, such that we get removable singularities. □

2.2 The Mellin transform

Definition: Let $f : \mathbb{R}^{\geq 0} \rightarrow \mathbb{R}$ be continuous. The Mellin transform $g(s)$ of f is defined by

$$g(s) := \int_0^\infty f(t) t^s \frac{dt}{t}$$

for values s such that the integral converges.

Example: The Mellin transform of e^{-t} is $\Gamma(s)$.

Example:

$$f(t) = e^{-ct} \Rightarrow g(s) = c^{-s} \Gamma(s) \quad (5)$$

2.3 Fourier transform

Definition: Let \mathcal{S} be the vector space of infinitely differentiable functions $f : \mathbb{R} \rightarrow \mathbb{C}$ such that $\lim_{x \rightarrow \pm\infty} |x|^n f(x) \rightarrow 0 \forall n \in \mathbb{N}$. For any $f \in \mathcal{S}$ we define the Fourier transform

$$\hat{f}(y) := \int_{-\infty}^{\infty} e^{-2\pi ixy} f(x) dx$$

and the integral converges for all $y \in \mathbb{C}$, $f \in \mathcal{S}$.

Example: Let $f(x) := e^{-\pi x^2}$, then $\hat{f} = f$.

Proof: Differentiating under the integral sign gives

$$\hat{f}'(y) = \frac{d}{dy} \int_{-\infty}^{\infty} e^{-2\pi ixy} f(x) dx = 2\pi i \int_{-\infty}^{\infty} e^{-2\pi ixy} x e^{-\pi x^2} dx$$

Integrating by parts yields

$$\begin{aligned} \hat{f}'(y) &= -2\pi i e^{-2\pi ixy} \frac{1}{-2\pi} e^{-\pi x^2} \Big|_{-\infty}^{\infty} + 2\pi i \int_{-\infty}^{\infty} 2\pi i y e^{-2\pi ixy} \frac{e^{-\pi x^2}}{-2\pi} dx \\ &= 2\pi y \int_{-\infty}^{\infty} e^{-2\pi ixy} f(x) dx = -2\pi y \hat{f}(y) \end{aligned}$$

Thus we have the differential equation

$$\frac{\hat{f}'(y)}{\hat{f}(y)} = -2\pi,$$

with the solution $\hat{f}(y) = C e^{-\pi y^2}$. Setting $y = 0$ gives

$$C = \hat{f}(0) = \int_{-\infty}^{\infty} e^{-\pi x^2} dx = 1$$

and thus $\hat{f}(y) = e^{-\pi y^2} = f(y)$.

Lemma: Let $f \in \mathcal{S}$ and $g(x) := f(ax)$ for some $a > 0$. Then $\hat{g}(y) = \frac{1}{a} \hat{f}(\frac{y}{a})$.

Proof:

$$\begin{aligned} \hat{g}(y) &= \int_{-\infty}^{\infty} e^{-2\pi ixy} f(ax) dx \\ &= \int_{-\infty}^{\infty} e^{-2\pi i \frac{x}{a} y} f(x) \frac{dx}{a} \\ &= \frac{1}{a} \hat{f}\left(\frac{y}{a}\right) \end{aligned}$$

□

Proposition (Poisson Summation): If $g \in \mathcal{S}$, then

$$\sum_{m=-\infty}^{\infty} g(m) = \sum_{m=-\infty}^{\infty} \hat{g}(m).$$

Proof: Define $h(x) := \sum_{k=-\infty}^{\infty} g(x+k)$. Clearly this has period 1. Write down the Fourier series:

$$h(x) = \sum_{m=-\infty}^{\infty} c_m e^{2\pi i m x}$$

with

$$\begin{aligned} c_m &:= \int_0^1 h(x) e^{-2\pi i m x} dx \\ &= \int_0^1 \sum_{k=-\infty}^{\infty} g(x+k) e^{-2\pi i m x} dx \\ &= \sum_{k=-\infty}^{\infty} \int_0^1 g(x+k) e^{-2\pi i m x} dx \\ &= \int_{-\infty}^{\infty} g(x) e^{-2\pi i m x} dx \\ &= \hat{g}(m) \end{aligned}$$

Then

$$\sum_{k=-\infty}^{\infty} g(k) = h(0) = \sum_{m=-\infty}^{\infty} c_m e^{-2\pi i m \cdot 0} = \sum_{m=-\infty}^{\infty} c_m = \sum_{m=-\infty}^{\infty} \hat{g}(m)$$

□

2.4 The theta function

Remark: Note that $\theta(t)$ can not only be written as in (2), but also as

$$\theta(t) = 1 + 2 \sum_{n=1}^{\infty} e^{-\pi n^2 t}$$

We will often make use of that fact.

Proposition 4: $\theta(t)$ satisfies the functional equation

$$\theta(t) = \frac{1}{\sqrt{t}} \theta(1/t)$$

Proof: Let $g(x) := e^{-\pi tx^2}$ for a fixed $t > 0$ and $f(x) := e^{-\pi x^2}$. Obviously $g(x) = f(\sqrt{t}x)$. Using the lemma and the example from the Fourier transform we get

$$\hat{g}(y) = \frac{1}{\sqrt{t}} \hat{f}(y/\sqrt{t}) = \frac{1}{\sqrt{t}} f(y/\sqrt{t}) = \frac{1}{\sqrt{t}} e^{-\pi y^2/t}$$

Using Poisson summation we finally find that

$$\begin{aligned} \theta(t) &:= \sum_{n=-\infty}^{\infty} e^{-\pi tn^2} = \sum_{n=-\infty}^{\infty} g(n) = \sum_{n=-\infty}^{\infty} \hat{g}(n) = \frac{1}{\sqrt{t}} \sum_{n=-\infty}^{\infty} e^{-\pi n^2/t} \\ &= \frac{1}{\sqrt{t}} \theta(1/t) \end{aligned}$$

□

Remark: In this proof we only need $\theta(t)$ for real-valued $t > 0$. But actually can also be looked at as a complex function for $\operatorname{Re} t > 0$ by analytic continuation. The functional equation still holds.

Proposition 5: As t goes to zero from above,

$$\left| \theta(t) - \frac{1}{\sqrt{t}} \right| < e^{-C/t} \quad \text{for some } C > 0.$$

Proof: Using Proposition 4 and a rewrite of θ gives

$$\left| \theta(t) - \frac{1}{\sqrt{t}} \right| = \left| \frac{1}{\sqrt{t}} (\theta(1/t) - 1) \right| = \frac{1}{\sqrt{t}} \cdot 2 \sum_{n=1}^{\infty} e^{-\pi n^2/t}$$

Suppose t is small enough such that $\sqrt{t} > 4 \cdot e^{-1/t}$ and $e^{-3\pi/t} < 1/2$. Then

$$\begin{aligned} \left| \theta(t) - \frac{1}{\sqrt{t}} \right| &< \frac{1}{2} e^{1/t} \left(e^{-\pi/t} + e^{-4\pi/t} + \dots \right) \\ &< \frac{1}{2} e^{-(\pi-1)/t} \left(1 + \frac{1}{2} + \frac{1}{4} + \dots \right) \\ &= e^{-(\pi-1)/t} \end{aligned}$$

and we see that $C = \pi - 1$ satisfies the inequality. □

3 The Theorem and its proof

Theorem $\zeta(s)$, as defined by (1) extends analytically onto the whole s -plane, except a simple pole at $s = 1$. Let

$$\Lambda(s) := \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s).$$

Then

$$\Lambda(s) = \Lambda(1 - s).$$

Proof: We want to consider the Mellin transform of the theta function,

$$\int_0^\infty \theta(t) t^s \frac{dt}{t}$$

As t goes to infinity, $\theta(t)$ converges rapidly against 1, as all terms of the sum fall to zero rapidly, except for $n = 0$. By proposition 5 we see that for small t , $\theta(t)$ behaves like $t^{-1/2}$. Thus, if we want convergence at both ends, we have to introduce correction terms and define

$$\phi(s) := \int_1^\infty (\theta(t) - 1) t^{s/2} \frac{dt}{t} + \int_0^1 \left(\theta(t) - \frac{1}{\sqrt{t}} \right) t^{s/2} \frac{dt}{t}.$$

Note that we replaced s by $s/2$ in order to get $\zeta(s)$ instead of $\zeta(2s)$. In the first integral, $\theta(t) - 1 \rightarrow 0$ extremely fast, such that the integral can be evaluated term by term, for any $s \in \mathbb{C}$. Similarly, $\theta(t) - \frac{1}{\sqrt{t}}$ is bounded above in the interval $(0, 1]$, such that the second integral converges for any $s \in \mathbb{C}$. These two statements together show that $\phi(s)$ is well-defined for all $s \in \mathbb{C}$, and it even is an entire function.

We now evaluate the second integral, assuming $\operatorname{Re} s > 1$:

$$\begin{aligned} & \int_0^1 \theta(t) t^{s/2} \frac{dt}{t} - \int_0^1 t^{(s-1)/2} \frac{dt}{t} \\ &= \int_0^1 \sum_{n=-\infty}^{\infty} e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} - \frac{2}{s-1} \\ &= \int_0^1 t^{s/2} \frac{dt}{t} + 2 \int_0^1 \sum_{n=1}^{\infty} e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} + \frac{2}{1-s} \\ &= 2 \int_0^1 \sum_{n=1}^{\infty} e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} + \frac{2}{s} + \frac{2}{1-s} \end{aligned}$$

Thus

$$\begin{aligned} \phi(s) &= 2 \int_1^\infty \sum_{n=1}^{\infty} e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} + 2 \int_0^1 \sum_{n=1}^{\infty} e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} + \frac{2}{s} + \frac{2}{1-s} \\ &= 2 \sum_{n=1}^{\infty} \int_0^\infty e^{-\pi n^2 t} t^{s/2} \frac{dt}{t} + \frac{2}{s} + \frac{2}{1-s} \end{aligned}$$

for $\operatorname{Re} s > 1$.

Using (5) we can evaluate the integral and we get

$$\begin{aligned} \frac{1}{2} \phi(s) &= \sum_{n=1}^{\infty} (\pi n^2)^{-s/2} \Gamma\left(\frac{s}{2}\right) + \frac{1}{s} + \frac{1}{1-s} \\ &= \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s) + \frac{1}{s} + \frac{1}{1-s} \\ \Rightarrow \zeta(s) &= \frac{\pi^{s/2}}{\Gamma\left(\frac{s}{2}\right)} \left(\frac{1}{2} \phi(s) - \frac{1}{s} - \frac{1}{1-s} \right) \end{aligned}$$

still for $\operatorname{Re} s > 1$. The only possible poles on the right hand side of the last equation are at $s = 0$ and $s = 1$, as $1/\Gamma(s)$ and $\phi(s)$ are entire functions. But for $s = 0$, the pole is removable, as the term causing the trouble is in fact

$$\frac{\pi^{s/2}}{\Gamma(\frac{s}{2})} \frac{1}{s} = \frac{\pi^{s/2}}{2 \cdot \frac{s}{2} \Gamma(\frac{s}{2})} = \frac{\pi^{s/2}}{2\Gamma(\frac{s}{2} + 1)} \xrightarrow{s \rightarrow 1} \frac{1}{2\Gamma\frac{3}{2}}.$$

Thus we have found a meromorphic function on \mathbb{C} with a (simple) pole at $s = 1$, and which equals $\zeta(s)$ for $\operatorname{Re} s > 1$. This is our analytic continuation. It remains to prove that the functional equation

$$\Lambda(s) = \Lambda(1 - s)$$

holds. Using

$$\frac{1}{2}\phi(s) = \pi^{-s/2}\Gamma(\frac{s}{2})\zeta(s) + \frac{1}{s} + \frac{1}{1-s} = \Lambda(s) + \frac{1}{s} + \frac{1}{s-1}$$

we have

$$\begin{aligned} \Lambda(s) &= \frac{1}{2}\phi(s) - \frac{1}{s} - \frac{1}{1-s} \\ \Lambda(1-s) &= \frac{1}{2}\phi(1-s) - \frac{1}{1-s} - \frac{1}{s} \end{aligned}$$

and we only have to prove $\phi(s) = \phi(1-s)$:

$$\begin{aligned} \phi(s) &= \int_1^\infty (\theta(t) - 1)t^{s/2} \frac{dt}{t} + \int_0^1 \left(\theta(t) - \frac{1}{\sqrt{t}} \right) t^{s/2} \frac{dt}{t} \\ &\stackrel{t \rightarrow \frac{1}{t}}{=} \int_0^1 \left(\theta\left(\frac{1}{t}\right) - 1 \right) t^{-s/2} \frac{dt}{t} + \int_1^\infty \left(\theta\left(\frac{1}{t}\right) - \sqrt{t} \right) t^{-s/2} \frac{dt}{t} \\ &= \int_0^1 (\sqrt{t}\theta(t) - 1)t^{-s/2} \frac{dt}{t} + \int_1^\infty (\sqrt{t}\theta(t) - \sqrt{t}) t^{-s/2} \frac{dt}{t} \\ &= \int_0^1 \left(\theta(t) - \frac{1}{\sqrt{t}} \right) t^{(1-s)/2} \frac{dt}{t} + \int_1^\infty (\theta(t) - 1) t^{(1-s)/2} \frac{dt}{t} \\ &= \phi(1-s) \end{aligned}$$

□