

Q1 Affine Connections

Let $\nabla: \Gamma(TM)^2 \rightarrow \Gamma(TM)$ be a given affine connection (so it satisfies the axioms on the bottom of pp. 18 in the lecture notes). (Recall $\Gamma(TM)$ is the $\mathcal{F}(M)$ -module of sections on TM , or of vector fields on M).

Let $\mathcal{S}: \Gamma(TM)^2 \rightarrow \Gamma(TM)$ be another given map (not necessarily an affine connection).

Define $\mathcal{B}: \Gamma(TM)^2 \rightarrow \Gamma(TM)$ via $\boxed{\mathcal{B} := \nabla - \mathcal{S}}$.

Note this makes sense as $\Gamma(TM)$ is an $\mathcal{F}(M)$ -module, so that its points may be added or subtracted.

Define a map $\mathcal{Q}: \Gamma(T^*M) \times \Gamma(TM)^2 \rightarrow \mathcal{F}(M)$ via
 $(w, X, Y) \mapsto w(\mathcal{B}(X, Y))$

Cl: \mathcal{S} is an affine connection $\iff \mathcal{Q} \in \Gamma(T^*M \otimes T^*M \otimes T^*M)$, that is, \mathcal{Q} is a tensor field of type (1,2).

Pr: Recall $\mathcal{Q} \in \Gamma(T^*M \otimes T^*M \otimes T^*M)$ iff it is an $\mathcal{F}(M)$ -multilinear map in each of its slots.

For \mathcal{Q} to be $\mathcal{F}(M)$ -linear in its 1st slot we should have:

① $\mathcal{Q}(fw, X, Y) \stackrel{?}{=} f \mathcal{Q}(w, X, Y) \quad \forall f \in \mathcal{F}(M)$.

But $\mathcal{Q}(fw, X, Y) \equiv (fw)(\mathcal{B}(X, Y))$
 $= f w(\mathcal{B}(X, Y))$
 $\equiv f \mathcal{Q}(\mathcal{B}(X, Y))$

② $\mathcal{Q}(w_1 + w_2, X, Y) \stackrel{?}{=} \mathcal{Q}(w_1, X, Y) + \mathcal{Q}(w_2, X, Y)$

But $\mathcal{Q}(w_1 + w_2, X, Y) \equiv (w_1 + w_2)(\mathcal{B}(X, Y))$
 $\equiv w_1(\mathcal{B}(X, Y)) + w_2(\mathcal{B}(X, Y))$
 $\equiv \mathcal{Q}(w_1, X, Y) + \mathcal{Q}(w_2, X, Y)$

So apparently \mathcal{Q} is always $\mathcal{F}(M)$ -linear in its first slot, regardless of what \mathcal{B} is. \checkmark

Next note that \mathcal{Q} is $\mathcal{F}(M)$ -linear in its 2nd slot iff \mathcal{B} is $\mathcal{F}(M)$ -linear in its 1st slot (by $\mathcal{F}(M)$ -linearity of w).

Since ∇ is $\mathcal{F}(M)$ -linear in its 1st slot and $\nabla = \mathcal{B} + \mathcal{S}$,

\mathcal{B} is $\mathcal{F}(M)$ -linear in its 1st slot iff \mathcal{S} is $\mathcal{F}(M)$ -linear in its 1st slot, which is the (i) and 1/2(ii) of the axioms for \mathcal{S} to be an affine connection.

2] Finally, Q is $\mathcal{F}(M)$ -linear in its 2nd slot iff B is $\mathcal{F}(M)$ -lin. in its 2nd slot. But note that

$$\nabla_x Y - \delta_x Y \equiv B(x, Y)$$

$$\Rightarrow \begin{cases} f \nabla_x Y - f \delta_x Y = f B(x, Y) \\ \nabla_x f Y - \delta_x f Y = B(x, f Y) \end{cases}$$

But we know that ∇ is an affine connection, so from its (iii)-axiom $\nabla_x f Y = (X(f))Y + f \nabla_x Y$

$$\Rightarrow \begin{cases} f \nabla_x Y - f \delta_x Y = f B(x, Y) \\ X(f)Y + f \nabla_x Y - \delta_x f Y = B(x, f Y) \end{cases}$$

$$\Leftrightarrow B(x, f Y) - f B(x, Y) = X(f)Y + f \nabla_x Y - \delta_x f Y$$

$$\Leftrightarrow \underbrace{\text{L.H.S.} = 0}_{B \text{ is } \mathcal{F}(M)\text{-lin. in 2nd slot}} \text{ iff } \underbrace{\text{R.H.S.} = 0}_{\delta \text{ obeys axiom (iii)}}$$

We find (separately for each axiom) that Q is $\mathcal{F}(M)$ -linear in its 2nd & 3rd slots iff δ is also an affine connection.

An equivalent claim:

Cl. The space of affine connections on M forms an affine space over the vector space of $\mathcal{F}(M)$ -bilinear maps $\Gamma(TM)^2 \rightarrow \Gamma(TM)$.

Pf. Recall A is an affine space over the v/sp. V iff \exists gp. morphism $t: V \rightarrow \text{Sym}(A)$ (where V is viewed as a gp. w.r.t. its additive structure and $\text{Sym}(A)$ is the gp. of bijections $A \xleftrightarrow{\sim} A$) s.t. $\forall a \in A$, $V \ni v \mapsto t(v)a \in A$ is bijective.

Let $A(M)$ be the sp. of affine connections on M , and $V(M)$ be the v/sp. of $\mathcal{F}(M)$ -bilinear maps $\Gamma(TM)^2 \rightarrow \Gamma(TM)$.

Define $t: V(M) \rightarrow \text{Aut}(A(M))$ via $B \mapsto (\cdot \mapsto \cdot - B)$

Verify it's a gp. morphism, that it's well-defined essentially by the proof above and it is bijective again by the same proof. (3)

Cl.: If $(\nabla, \delta) \in \mathcal{A}(\mathcal{M})^2$ then $((1-\alpha)\nabla + \alpha\delta) \in \mathcal{A}(\mathcal{M})$
 $\forall \alpha \in (0,1)$.

Pf.: Let B be as before w.r.t. ∇ and δ .

$$\text{Nde } \nabla - [(1-\alpha)\nabla + \alpha\delta] = \alpha(\nabla - \delta) \equiv \alpha B$$

Since $B \in \mathcal{V}(\mathcal{M}) \Leftrightarrow \alpha B \in \mathcal{V}(\mathcal{M})$ (with $\mathcal{V}(\mathcal{M})$ as above) $\forall \alpha \in \mathbb{R}$, we apply the 1st claim to get the result.

Let $\varphi: \mathcal{U} \rightarrow \mathbb{R}^n$ be a chart. Recall we found an expression for the Christoffel symbols w.r.t. φ of the parallel transport induced by a $\nabla \in \mathcal{A}(\mathcal{M})$ to be:

$$\Gamma_{ijk}^\varphi = e_i^\varphi \left(\nabla_{d_j^\varphi} d_k^\varphi \right)$$

where $\{e_i^\varphi\}_i$ was the chart-basis for $\Gamma(T^*\mathcal{M}|_{\mathcal{U}})$, and $\{d_i^\varphi\}_i$ $\Gamma(T\mathcal{M}|_{\mathcal{U}})$.

Cl.: The components of Q , Q^φ_{ijk} , satisfy:

$$Q^\varphi_{ijk} = \Gamma_{ijk}^\varphi - \tilde{\Gamma}_{ijk}^\varphi$$

$$\text{Pf.} \quad \Gamma_{ijk}^\varphi - \tilde{\Gamma}_{ijk}^\varphi = e_i^\varphi \left(\nabla_{d_j^\varphi} d_k^\varphi \right) - e_i^\varphi \left(\tilde{\nabla}_{d_j^\varphi} d_k^\varphi \right)$$

$$\stackrel{e_i^\varphi \text{ linear}}{\cong} e_i^\varphi \left(\nabla_{d_j^\varphi} d_k^\varphi - \tilde{\nabla}_{d_j^\varphi} d_k^\varphi \right)$$

$$\stackrel{\text{def. of } B}{\cong} e_i^\varphi \left(B(d_j^\varphi, d_k^\varphi) \right)$$

$$\cong Q(e_i^\varphi, d_j^\varphi, d_k^\varphi) \equiv Q^\varphi_{ijk}$$

We see again that the difference of two Christoffel symbols of two different connections behaves like a tensor, just like the difference of the symbols in two different charts.

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Q2

Euclidean Metric in Polar Coord.

$M = \mathbb{R}^2$ with the chart $\varphi = \mathbb{1}_M$.

The metric g is a point in $\Gamma(T^*M \otimes T^*M)$ which obeys certain axioms.

Thus we may specify it as the components of a (0,2)-tensor in a chart-basis $\{e_i^\varphi\}$ as:

$$\{g(d_i^\varphi, d_j^\varphi)\}_{i,j=1}^2$$

So, a 2×2 matrix.

The axioms are that g should be symmetric & non-deg., for that components matrix that means it should be symmetric and invertible.

Define $\{g(d_i^\varphi, d_j^\varphi)\}_{i,j} := \mathbb{1}_{2 \times 2}$, which satisfies these requirements.

Recall from the previous HW the polar coord. chart $\varphi: \mathbb{R}^2 \setminus \{0\} \rightarrow (0, \infty) \times [0, 2\pi)$

$$x \mapsto \begin{bmatrix} \|x\| \\ \arctan(x_2/x_1) \end{bmatrix}$$

To compute the metric in polar coord. we apply the usual coord. transf. \circ

$$\underbrace{g(d_i^\varphi, d_j^\varphi)}_{g_{ij}^\varphi} = N_{ii}^{\varphi\varphi} N_{jj}^{\varphi\varphi} \underbrace{g(d_i^{\varphi\varphi}, d_j^{\varphi\varphi})}_{=\delta_{i'j'}} = N_{ij}^{\varphi\varphi} N_{j'i'}^{\varphi\varphi} \quad (|A|^2 = A^T A)$$

$$= (N^{\varphi\varphi} (N^{\varphi\varphi})^T)_{ij} = |(N^{\varphi\varphi})^T|_{ij}^2$$

We found then

$$N^{\varphi\varphi}(\varphi^{-1}(r, \varphi)) = \begin{bmatrix} \cos(\varphi) & \sin(\varphi) \\ -r\sin(\varphi) & r\cos(\varphi) \end{bmatrix}$$

$$\Rightarrow g^\varphi(\varphi^{-1}(r, \varphi)) = \begin{bmatrix} \swarrow & \downarrow \end{bmatrix} \begin{bmatrix} \downarrow \\ \swarrow \end{bmatrix}^T = \begin{bmatrix} 1 & 0 \\ 0 & r^2 \end{bmatrix}$$

The Christoffel symbols of the g-Levi-Civita connection (5)
are given by the formula

$$\Gamma_{ijk}^p = \frac{1}{2} (g^{pq})^{-1} (g_{jk,p} + g_{kp,j} - g_{jk,p})$$

Note $g_{,r}^k = \begin{bmatrix} 0 & 0 \\ 0 & 2r \end{bmatrix}$ $g_{,p}^k = 0$

$$(g^k)^{-1} = \begin{bmatrix} 1 & 0 \\ 0 & r^{-2} \end{bmatrix}$$

$$(g^k)^{-1} g_{,p}^k = 0$$

$$(g^k)^{-1} g_{,r}^k = \begin{bmatrix} 0 & 0 \\ 0 & 2r^{-1} \end{bmatrix}$$

$$\Gamma_{ijk}^k = \frac{1}{2} (g^k)^{-1}_{ir} (g_{,r}^k)_{jk}$$

Q3 Torsion of the Hessian

$\nabla \in \mathcal{A}(M)$ affine connection

Cl.: $df = \nabla f \quad \forall f \in \mathcal{F}(M)$

Pr.: $f: M \rightarrow \mathbb{R} \Rightarrow df: T_p M \rightarrow T_{p(p)} \mathbb{R} \cong \mathbb{R}$
 $X \mapsto X(\cdot \circ f) \equiv X(f)$

O.T.O.H., $\nabla \cdot f$ is a 1-form:

$(\nabla \cdot f)(X) \equiv \nabla_X f = X(f)$

Def.: $\forall f \in \mathcal{F}(M)$, the Hessian H of f is
 $H(p) := \nabla(\nabla f) \in \Gamma(T^*M \otimes T^*M)$.

Recall the def. of torsion: (2.8)

Def.: Torsion T corresp. to $\nabla \in \mathcal{A}(M)$ is a map
 $\Gamma(TM)^2 \rightarrow \Gamma(TM)$ defined by
 $T(X, Y) := \nabla_X Y - \nabla_Y X - [X, Y]$

Cl.: $H(p)$ is symmetric $\Leftrightarrow T = 0$

Pr.: $(H(f))(X, Y) \equiv (\nabla(df))(X, Y)$
 $\equiv (\nabla_Y df)(X) \stackrel{(2.8)}{=} Y(X(f)) - (df)(\nabla_Y X)$
 (Note: ∇ slot of ∇ always takes last argument)

$= (YX - \nabla_Y X) f$

$\Rightarrow (H(f))(X, Y) - (H(f))(Y, X) = (YX - \nabla_Y X - XY + \nabla_X Y) f$
 $= (\nabla_X Y - \nabla_Y X - [X, Y]) f$
 $\equiv (T(X, Y))(f)$

Q4

Geodesics in the Hyperbolic Plane

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$$M := \{x \in \mathbb{R}^2 \mid \pi_2(x) > 0\}$$

Chart $\varphi := \mathbb{1}|_M$ $\xrightarrow{\pi_2: \mathbb{R}^2 \rightarrow \mathbb{R}}$ proj. onto new coord.

$$\text{Metric } g^\varphi(x) = \pi_2(x)^{-2} \mathbb{1}_{2 \times 2} \equiv (x_2)^{-2} \mathbb{1}_{2 \times 2}$$

Recall that a curve $\gamma: \mathbb{R} \rightarrow M$ is called a g -geodesic iff $\boxed{\nabla \dot{\gamma} = 0}$ where ∇ is the g -Levi-Civita connection.

That implies in a chart $\nabla \dot{\gamma} = 0$ becomes (3.10):

$$\boxed{\ddot{\gamma}_i^\varphi + \Gamma_{ijk}^\varphi \dot{\gamma}_j^\varphi \dot{\gamma}_k^\varphi = 0} \quad \Gamma \text{ evaluated on } \dot{\gamma}!$$

We start by calculating Γ^φ :

$$\Gamma_{ijk}^\varphi = \frac{1}{2} (g^\varphi)^{-1}_{il} \left[(g^\varphi)_{je,k} + (g^\varphi)_{ke,j} - (g^\varphi)_{jk,e} \right]$$

$$g^\varphi(x) = (x_2)^{-2} \mathbb{1}_{2 \times 2}, \quad (g^\varphi)^{-1}(x) = (x_2)^2 \mathbb{1}_{2 \times 2}$$

$$\Gamma_{1jk}^\varphi(x) = \frac{1}{2} (x_2)^2 \left[\underbrace{(g^\varphi)_{j1,k}}_{\substack{[(x_2)^{-2} \delta_{j,1}]_k \\ -2(x_2)^{-3} \delta_{j,1} \delta_{k,2}}} + (g^\varphi)_{k1,j} - \underbrace{(g^\varphi)_{jk,1}}_0 \right]$$

$$= -\frac{1}{x_2} (\delta_{j,1} \delta_{k,2} + \delta_{j,2} \delta_{k,1})$$

$$\boxed{\Gamma_{1..}^\varphi(x) = -\frac{1}{x_2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}$$

$$\Gamma_{2jk}^\varphi(x) = \frac{1}{2} (x_2)^2 \left[\underbrace{(g^\varphi)_{j2,k}}_{\substack{[(x_2)^{-2} \delta_{j,2}]_k \\ -2(x_2)^{-3} \delta_{j,2} \delta_{k,2}}} + (g^\varphi)_{k2,j} - \underbrace{(g^\varphi)_{jk,2}}_{-2(x_2)^{-3} [\delta_{j,2} \delta_{k,1} + \delta_{j,2} \delta_{k,2}]} \right]$$

$$= \frac{1}{x_2} \delta_{j,1} \delta_{k,1} - \frac{1}{x_2} \delta_{j,2} \delta_{k,2} \Rightarrow \boxed{\Gamma_{2..}^\varphi(x) = \frac{1}{x_2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}}$$

We go on to calculate the geodesic eq-n:

$$\ddot{\gamma}_1^{\varphi} + \begin{bmatrix} \dot{\gamma}_1^{\varphi} & \dot{\gamma}_2^{\varphi} \end{bmatrix} \begin{bmatrix} 0 & -\frac{1}{\gamma_2^{\varphi}} \\ -\frac{1}{\gamma_2^{\varphi}} & 0 \end{bmatrix} \begin{bmatrix} \dot{\gamma}_1^{\varphi} \\ \dot{\gamma}_2^{\varphi} \end{bmatrix} = 0$$

$$\boxed{\gamma_2^{\varphi} \ddot{\gamma}_1^{\varphi} = 2 \dot{\gamma}_1^{\varphi} \dot{\gamma}_2^{\varphi}} \quad (1)$$

(We omit the φ superscript...)

$$\ddot{\gamma}_2 + \begin{bmatrix} \dot{\gamma}_1 & \dot{\gamma}_2 \end{bmatrix} \begin{bmatrix} \frac{1}{\gamma_2} & 0 \\ 0 & -\frac{1}{\gamma_2} \end{bmatrix} \begin{bmatrix} \dot{\gamma}_1 \\ \dot{\gamma}_2 \end{bmatrix} = 0$$

$$\boxed{\gamma_2 \ddot{\gamma}_2 = (\dot{\gamma}_2)^2 - (\dot{\gamma}_1)^2} \quad (2)$$

Cl.: If $\dot{\gamma}_1 \neq 0$ then $(1) \wedge (2) \Rightarrow \gamma_1^2 + \gamma_2^2 - a\gamma_1 = b$
 $\exists (a, b) \in \mathbb{R}^2$.

Pf.: $\dot{\gamma}_1 \neq 0 \Rightarrow (1)$ implies $\frac{\ddot{\gamma}_1}{\dot{\gamma}_1} = 2 \frac{\dot{\gamma}_2}{\gamma_2}$

Note: $\log(f) = \frac{1}{f} \dot{f} \Rightarrow \log(\dot{\gamma}_1) = 2 \log(\dot{\gamma}_2)$

$$\Rightarrow \log(\dot{\gamma}_1) = 2 \log(\dot{\gamma}_2) + c \quad \exists c \in \mathbb{R}$$

$$\Rightarrow \dot{\gamma}_1 = \exp(\log(\dot{\gamma}_2^2) + c) = \dot{\gamma}_2^2 \exp(c)$$

$$\boxed{\dot{\gamma}_1 = D \dot{\gamma}_2^2} \quad D := \exp(c)$$

Let $f: \mathbb{R} \rightarrow \mathbb{R}_{>0}$ be given by $\gamma_1(t) \mapsto \gamma_2(t)$
 (We know we can find one such map for all t by reparametrizing γ_2 to be wrt. γ_1 instead of t)

Cl.: (2) implies $f f'' + (1 + f')^2 = 0$

$$\text{Pf.} \quad f' = \frac{\partial \gamma_2}{\partial \gamma_1} = \frac{\dot{\gamma}_2}{\dot{\gamma}_1}$$

$$f'' = \left(\frac{\dot{\gamma}_2}{\dot{\gamma}_1} \right)' = \frac{\ddot{\gamma}_2}{(\dot{\gamma}_1)^2} - \frac{\dot{\gamma}_2}{(\dot{\gamma}_1)^2} (\dot{\gamma}_1)'$$

$$(\dot{\gamma}_1)' = (Df^2)' = 2f f' = 2D\gamma_2 \frac{\dot{\gamma}_2}{\dot{\gamma}_1} = \frac{2\dot{\gamma}_2}{\dot{\gamma}_1} \quad (7)$$

$$\begin{aligned} \Rightarrow f f'' + 1 + f'^2 &= \frac{\gamma_2 \ddot{\gamma}_2}{\dot{\gamma}_1^2} - \frac{2\dot{\gamma}_2^2}{\dot{\gamma}_1^3} + 1 + \frac{\dot{\gamma}_2^2}{\dot{\gamma}_1^2} = \\ &= \dot{\gamma}_1^{-2} (\gamma_2 \ddot{\gamma}_2 - \dot{\gamma}_2^2 + \dot{\gamma}_1^2) \stackrel{\text{by (2)}}{=} 0 \end{aligned}$$

However $f f'' + 1 + f'^2 = 0$ may be solved via:

$$f(x) = \pm \sqrt{a^2 - (x-x_0)^2} \quad (\text{the - solution is irrelevant})$$

Indeed, $f'(x) = \frac{1}{2} f(x)^{-1} (-2(x-x_0)) = -\frac{x-x_0}{f(x)}$

$$f''(x) = -\frac{1}{f(x)} + \frac{x-x_0}{f(x)^2} f'(x) = \frac{1}{f(x)} - \frac{(x-x_0)^2}{f(x)^3}$$

$$(f f'' + 1 + f'^2)(x) = \cancel{-1} - \frac{(x-x_0)^2}{f(x)^2} + \cancel{1} + \frac{(x-x_0)^2}{f(x)^2} = 0 \quad \checkmark$$

We find the desired result.

This eqn describes a semi-circle whose center is on the horizontal axis.

Note that if $\dot{\gamma}_1 = 0$, we find γ describes a vertical line.

Note that except for $\frac{\dot{\gamma}_1}{\dot{\gamma}_2}$ being a constant there

is another constant of motion:

Q.: $\frac{\|\dot{\gamma}\|^2}{\dot{\gamma}_2^2}$ is a constant of motion.

PP.: The geodesic γ is the extremum of the Lagrangian

$$L(\gamma_1, \gamma_2, \dot{\gamma}_1, \dot{\gamma}_2) = \frac{1}{2} \frac{\|\dot{\gamma}\|^2}{\dot{\gamma}_2^2}$$

Since L doesn't depend explicitly on time, its Legendre transform, the "energy", is const.

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$$\dot{q}_2 \frac{\partial L}{\partial \dot{q}_2} + \dot{q}_1 \frac{\partial L}{\partial \dot{q}_1} - L = L$$

Later on you'll see another way to find conserved quantities via "Killing fields".