

(Q1) (Alternative way to the one suggested by the hints)

In (Q2) we find that for a grav. field φ ,

$$\Gamma_{i00} = \partial_i \varphi \quad \forall i \in \{1, 2, 3\}$$

Work in 2D spacetime. $\Rightarrow \Gamma_{100} = \varphi'$

Assume metric doesn't dep. on time:

$$\begin{aligned} \Gamma_{100} &= \frac{1}{2} (g^{-1})_{1\alpha} (g_{\alpha e, 0} + g_{0e, \alpha} - g_{00, e}) \\ &= \frac{1}{2} (g^{-1})_{1\alpha} g_{00, \alpha} \\ &= -\frac{1}{2} (g^{-1})_{11} g_{00, 1} \end{aligned}$$

Assume g is almost the Minkowski metric: $g = \begin{bmatrix} +1 & 0 \\ 0 & -1 \end{bmatrix} + h$

where h has "small" components.

$$\Rightarrow g^{-1} \approx \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + O(h), \quad g_{00, 1} = h_{00, 1} = h_{00}'$$

$$\Rightarrow \boxed{\Gamma_{100} = +\frac{1}{2} h_{00, 1}} \Rightarrow \frac{1}{2} h_{00}' = \varphi'$$

$$\Rightarrow \boxed{g_{00} = 1 + 2\varphi} \quad (\text{for this choice of const})$$

If we pick a homogeneous grav. field, $\varphi(x) = \gamma x_1$ where γ is the gravitational const. (9.8) with γ instead of g to avoid conflict of notation with the metric g .

We know (eqn (4.4) in the script) that

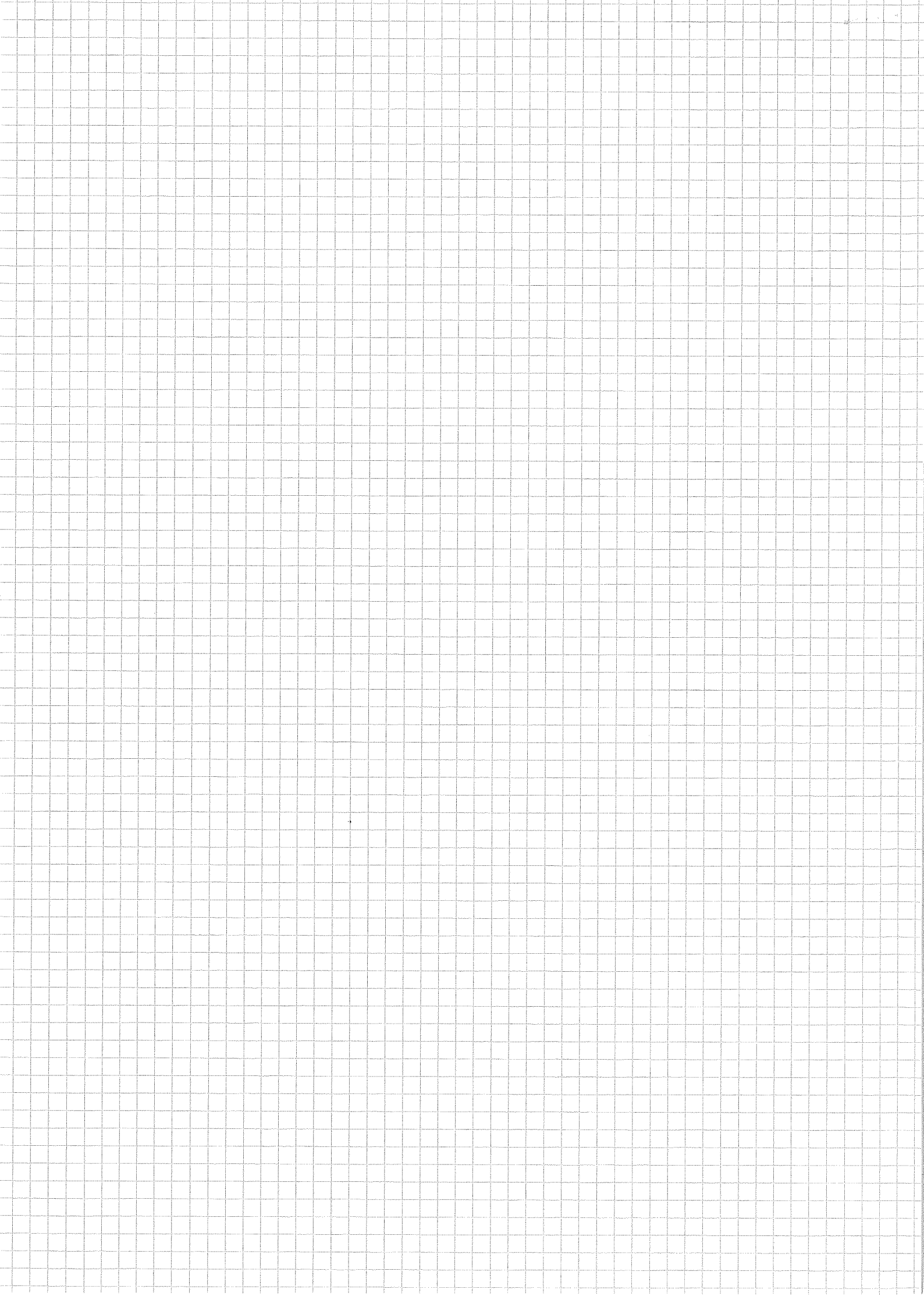
$$(\Delta \tau)^2 = g_{00}(x) (\Delta t)^2 \Rightarrow \boxed{\Delta \tau = \sqrt{1 + 2\gamma x_1} \Delta t}$$

where $\Delta \tau$ is the ^{elapsed} proper time of the clock, x_1 is its position, and Δt is the elapsed time in the chart.

We have one clock at height l (φ_2) and another clock at height zero (φ_1) $\Rightarrow \Delta \tau_1 = \Delta t$, $\Delta \tau_2 = \sqrt{1 + 2\gamma l} \Delta t$.

$$\Rightarrow \Delta \tau_1 = (1 + 2\gamma l)^{-1/2} \Delta \tau_2 \approx (1 - \gamma l) \Delta \tau_2$$

\uparrow
 $\gamma l \ll 1$



Q2. Newton's Eq-n as a Geodesic Eq-n 3

In classical mechanics, when a particle is placed in a gravitational potential $\varphi: \mathbb{R}^3 \rightarrow \mathbb{R}$ its trajectory $\gamma: \mathbb{R} \rightarrow \mathbb{R}^3$ obeys the equation $\ddot{\gamma} = -(\nabla\varphi) \circ \gamma$. \otimes

In 4D spacetime, a curve is given by a map $\tilde{\eta}: \mathbb{R} \rightarrow \mathcal{M}$ where the domain \mathbb{R} is the parameter of the curve $\tilde{\eta}$ (not necessarily time) and \mathcal{M} is the spacetime manifold, which by GR is a 4D pseudo-Riemannian manifold.

Hence we should have a chart $\psi: \mathcal{U} \rightarrow \mathbb{R}^4 \exists \mathcal{U} \in \text{Open}(\mathcal{M})$.
Then $\psi \circ \tilde{\eta}|_{\tilde{\eta}^{-1}(\mathcal{U})}: \tilde{\eta}^{-1}(\mathcal{U}) \rightarrow \mathbb{R}^4$ is a curve in \mathbb{R}^4 parametrized by an open subset of \mathbb{R} (since $\tilde{\eta}$ is cont.).

$$V := \tilde{\eta}^{-1}(\mathcal{U}) \in \text{Open}(\mathbb{R})$$

$$\eta := \psi \circ \tilde{\eta}|_V$$

η is hence a curve in \mathbb{R}^4 .

If we want to "connect" η with γ then we could "declare" that $\eta_0 := \downarrow$ (viewing the 0th component of \mathbb{R}^4 as that which corresponds to time) and $(\eta_i)_{i=1,2,3} := \gamma$.

Then we can figure out which EoM η satisfies:

We let now the dot be the derivative w.r.t. the param. of the curve η , $\lambda \in V$:

$$\dot{\eta}_0(\lambda) \equiv 1 \Rightarrow \ddot{\eta}_0(\lambda) = 0$$

From \otimes we get:

$$\begin{aligned} (\partial_t^2 \gamma)_i(t) &= -(\nabla\varphi)_i(\gamma(t)) = -(\partial_i \varphi)(\gamma(t)) \\ \Rightarrow \ddot{\eta}_i &= (\partial_\lambda^2 \eta)_i(\lambda) + (\partial_i \varphi)(\eta(\lambda)) = 0 \end{aligned}$$

But $\dot{\eta}_0(\lambda) = 1$, so we may rewrite this as:

$$\ddot{\eta}_i + [(\partial_i \varphi) \circ \eta] \dot{\eta}_0 \dot{\eta}_0 = 0$$

We then recall what the geodesic eq-n for a curve $\eta: V \rightarrow \mathbb{R}^4$ would look like given Christoffel symbols:

$$\ddot{\eta}_i + \Gamma_{ijk} \dot{\eta}_j \dot{\eta}_k = 0$$

For the two equations to coincide, we may define

$$\Gamma_{0jk} := 0 \quad \text{for } \dot{\eta}_0 = 0.$$

$$\Gamma_{i00} := 2\delta_{i4} \quad \text{for the other three eqns, } i \in \{1, 2, 3\}.$$

All other components should be zero.

Such a def. for Γ is clearly symmetric, giving (i). ✓

The corresponding Riemann curvature tensor, given by eqn (2.15) in the script:

$$R_{ijkl} \stackrel{(2.15)}{=} \Gamma_{ilj,k} - \Gamma_{ikj,l} + \Gamma_{slj}\Gamma_{iks} - \Gamma_{skj}\Gamma_{ils} \quad (i,j,k,l) \in \{0,1,2,3\}$$

$$\begin{aligned} \Gamma_{slj}\Gamma_{iks} - \Gamma_{skj}\Gamma_{ils} &= \underbrace{\Gamma_{0lj}\Gamma_{iko}}_{=0} + \sum_{s=1}^3 \Gamma_{slj}\Gamma_{iks} - \underbrace{\Gamma_{0kj}\Gamma_{ilo}}_{=0} - \sum_{s=1}^3 \Gamma_{skj}\Gamma_{ils} \\ &= -\sum_{s=1}^3 \Gamma_{skj}\Gamma_{ils} = 0 \quad \text{if } k \in \{1, 2, 3\}. \end{aligned}$$

Also $\Gamma_{ikj,l} = 0$ if $k \in \{1, 2, 3\}$.

$$\Rightarrow R_{ijkl} = \Gamma_{ilj,k} \quad \text{if } k \in \{1, 2, 3\}.$$

$$\Rightarrow R_{i0k0} = \Gamma_{i00,k} = 2k\delta_{i4} \quad \text{if } k \in \{1, 2, 3\}.$$

$$R_{ij00} = \Gamma_{i0s,k} = 0 \quad \text{if } i,j,k \in \{1, 2, 3\}.$$

If Γ were to be associated with a Levi-Civita connection, we know it should obey the symmetries on pp. 29, namely:

$$\begin{cases} g_{ii} R_{ijkl} = -g_{jj} R_{jikl} \\ g_{ii} R_{ijkl} = g_{kk} R_{klij} \end{cases}$$

$$g_{ii} R_{i0k0} = g_{ii} 2k\delta_{i4} \stackrel{?}{=} -g_{0j} \underbrace{R_{j0k0}}_{=0}$$

\Rightarrow Not true unless $2k\delta_{i4} = 0 \Leftrightarrow$ Const. grav. field.

B3. On the Levi-Civita Connection

Let M be a manifold.

Def.: A subset $N \subseteq M$ is called an embedded submanifold of M iff $\forall p \in N \exists$ chart $\psi: U \rightarrow \mathbb{R}^m$ ($U \in \text{Open}(M)$, $p \in U$) st.
 $\psi(U \cap N) = \psi(U) \cap W$ where $W \subseteq \mathbb{R}^m$ is a subspace of dim. $n = \dim(N)$, $0 \leq n \leq m$. The atlas for N is the pairs $(U \cap N, \psi|_{U \cap N})$.

$\iota: N \hookrightarrow M$, the inclusion map, induces a tangent map:

$(d\iota)_p: T_p N \hookrightarrow T_p M$, which is injective as well. (We also use $\iota_* \equiv d\iota$)

$$(d\iota)_p(\iota_* v) = (d\iota)_p(\iota_* v)$$

$$\begin{aligned} \updownarrow \\ \iota_* (\iota_* v) &= \iota_* (\iota_* v) \\ \updownarrow \end{aligned}$$

$$\iota_* (\iota_* v) = \iota_* (\iota_* v) \quad \forall v \in T_p N$$

But since $\forall g \in \mathcal{F}(N)$, \exists extension $\tilde{g} \in \mathcal{F}(M)$: $\tilde{g} \circ \iota = g$, $\iota_* = \iota_*$.

\Rightarrow We think of $T_p N$ as a linear subsp. of $T_p M$.

If $g \in \Gamma(T^*M \otimes T^*M)$ is a metric on M , a \lfloor metric $h \in \Gamma(T^*N \otimes T^*N)$ is induced via the eq'n:

$$h(X, Y) := g(d\iota(X), d\iota(Y))$$

$$\forall (X, Y) \in \Gamma(TN)^2.$$

$\forall p \in N$, define $P_p: T_p M \rightarrow T_p N$ as the orthogonal proj. assoc. to g : g is an inner-prod. on $T_p M$, so it defines what it means for two vectors to be orthogonal:

$$[u \perp_p v \iff g_p(u, v) = 0] \quad \forall (u, v) \in (T_p M)^2$$

Then the orthogonal projection P_p satisfies:

$$\begin{aligned} P_p u &= u \quad \forall u \in T_p N \quad \text{bec. } P_p \text{ is a proj.} \\ g_p(v, d\iota u) &= g_p(v, P_p(d\iota u)) = g_p(P_p v, d\iota u) \\ &\quad \uparrow \\ &\text{bec. it is an orthogonal proj.} \end{aligned}$$

$$\forall u \in T_p N, v \in T_p M.$$

Let ∇ be the g -Levi-Civita connection.

Cl. 0 $(\nabla_X Y)_p = \tilde{P}_p(\nabla_{\tau_X} \tau_* Y) \quad \forall (X, Y) \in \Gamma(TN)^2, p \in \mathcal{N}$.
 with \tilde{P} being the proj. \mathbb{R} -l. w/ its codomain TN .

Pf. 0 Let $(X, Y) \in \Gamma(TN)^2, p \in \mathcal{N}$.

Cl. 1 $\tilde{P} \circ \nabla_{\tau_X} \tau_*$ is an affine connection on \mathcal{N} .

Pf. 1 (i) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_*$ is $\mathcal{F}(N)$ -lin. in X :

$$\tilde{P} \circ \nabla_{\tau_{X_1+X_2}} \tau_* Y = \tilde{P} \circ \nabla_{\tau_{X_1} + \tau_{X_2}} \tau_* Y$$

\uparrow
 τ_* linear

∇ is an affine conn. $\Rightarrow \tilde{P} \circ (\nabla_{\tau_{X_1}} \tau_* Y + \nabla_{\tau_{X_2}} \tau_* Y)$

\tilde{P} is linear $\Rightarrow \tilde{P} \circ \nabla_{\tau_{X_1}} \tau_* Y + \tilde{P} \circ \nabla_{\tau_{X_2}} \tau_* Y \quad \checkmark$

$$(\tilde{P} \circ \nabla_{\tau_f X} \tau_* Y)_p = (\tilde{P} \circ \nabla_{f \tau_X} \tau_* Y)_p \quad \begin{matrix} f \in \mathcal{F}(N) \\ p \in \mathcal{N} \end{matrix}$$

$\tau_* X = f \tau_X$

$$= (\tilde{P} \circ f \nabla_{\tau_X} \tau_* Y)_p$$

$$= \tilde{P}_p(f(p) \nabla_{\tau_X} \tau_* Y)_p$$

\tilde{P}_p is linear $\Rightarrow f(p) \tilde{P}_p(\nabla_{\tau_X} \tau_* Y)_p \quad \checkmark$

(ii) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_*$ is \mathbb{R} -lin. in Y :

$Y_1 + Y_2$ follows by lin. of τ_* , ∇ being affine connection and \tilde{P} being linear.

(iii) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_*$ obeys Leibniz in Y :

$$(\tilde{P} \circ \nabla_{\tau_X} \tau_* fY)_p = \tilde{P}_p(\nabla_{\tau_X} (f \tau_* Y))_p =$$

$\tilde{P}_p \circ (f(p) \nabla_{\tau_X} \tau_* Y + \tau_* Y (\tau_X X)(f))_p$
 \tilde{P} extends f to \mathcal{M} (trivially)

$$= f(p) \tilde{P}_p(\nabla_{\tau_* X} \tau_* Y)_p + (\tilde{P}_p(\tau_* Y)_p)(\tau_* X)(f)$$

$$\tilde{P}_p(\tau_* \cdot)_p = 1 \Rightarrow \cong f(p) \tilde{P}_p(\nabla_{\tau_* X} \tau_* Y)_p + (Y)_p (\tau_* X)(f)$$

$\Rightarrow \tilde{P}_0 \nabla_{\tau_*} \tau_* \cdot$ is an affine conn. on \mathcal{N} . ✓

Cl: The torsion of $\tilde{P}_0 \nabla_{\tau_*} \tau_* \cdot$ is zero.

Pr: The torsion is given by:

$$\tilde{P}_0 \nabla_{\tau_* X} \tau_* Y - \tilde{P}_0 \nabla_{\tau_* Y} \tau_* X - [\tau_* X, \tau_* Y] =$$

$$\stackrel{\tilde{P}_0 \tau_* = 1}{=} \tilde{P}_0 (\nabla_{\tau_* X} \tau_* Y - \nabla_{\tau_* Y} \tau_* X - \tau_* [X, Y])$$

Cl: $\tau_* [X, Y] = [\tau_* X, \tau_* Y]$

Pr: $(\tau_* [X, Y])(p) \equiv ([X, Y])(\tau_0 p)$
 $= X(Y(\tau_0 p)) - Y(X(\tau_0 p))$
 $([\tau_* X, \tau_* Y])(p) \equiv (\tau_* X)(Y(\tau_0 p)) - \dots$
 $\equiv X(Y(\tau_0 p)) - \dots$
 $\stackrel{\tau_0(p) \equiv p}{\forall p \in \mathcal{N}}}{=} X(Y(\tau_0 p))$

$\Rightarrow = \tilde{P}_0 0$ as ∇ has no torsion.

Cl: $(\tilde{P}_0 \nabla_{\tau_*}) (h) = 0$

Pr: We use the inductive formula: $Z(T(X_1, \dots, X_k, \mu_1, \dots, \mu_l))$

$$(\nabla_Z T)(X_1, \dots, X_k, \mu_1, \dots, \mu_l) = \nabla_Z (T(X_1, \dots, X_k, \mu_1, \dots, \mu_l))$$

$$- T(\nabla_Z X_1, X_2, \dots, X_k, \mu_1, \dots, \mu_l)$$

$$\dots$$

$$- T(X_1, \dots, X_k, \mu_1, \dots, \nabla_Z \mu_l)$$

to get:
 $(\tilde{P}_0 \nabla_{\tau_*} h)(X, Y) = Z(h(X, Y)) - h(\tilde{P}_0 \nabla_{\tau_*} X, Y) - h(X, \tilde{P}_0 \nabla_{\tau_*} Y)$

$$= \underbrace{1_* Z}_{\text{bcs. } 1(p) \equiv p \ \forall p \in V} (g(1_* X, 1_* Y)) - \underbrace{g(1_* \tilde{P} \nabla_{1_* Z} 1_* X, 1_* Y)}_{P} - \underbrace{g(1_* X, 1_* \tilde{P} \nabla_{1_* Z} 1_* Y)}_P$$

P orthogonal
proj

$$= 1_* Z (g(1_* X, 1_* Y)) - g(\nabla_{1_* Z} 1_* X, \underbrace{P 1_* Y}_{= 1_* Y}) - g(\underbrace{P 1_* X}_{= 1_* X}, \nabla_{1_* Z} 1_* Y)$$

$$= (\nabla_{1_* Z} g)(1_* X, 1_* Y)$$

$$= 0 \quad \text{in particular,}$$

Hence by uniqueness of the Levi-Civita connection we obtain the result.