

Higgs-Coulomb correspondence in GLSM

GLSM is a curve counting theory (GW, quasimap, DT, ...)

Consider the following problem: let's count maps from

$$\begin{array}{ccc}
 \mathbb{P}^1 & \rightarrow & \mathbb{P}^1 \\
 \uparrow \mathbb{C}^* & & \uparrow \mathbb{C}^* \\
 (x:y) & & (X:Y)
 \end{array}
 \quad \rightarrow \quad
 \begin{array}{ccc}
 (x:y) & \rightarrow & (q x:y) \\
 (X:Y) & \rightarrow & (a_1 X : a_2 Y)
 \end{array}$$

Degree d map $(P(x,y) : Q(x,y))$, where P, Q are homog-s of degree d such that they do not have common zeros.

Let's compactify space of such maps.

- stable maps (GW theory)
- quasimaps (Ciocan-Fortunier, Kim)

Quasimap $\mathbb{P}^1 \dashrightarrow \mathbb{P}^1$ is $(P(x,y) ; Q(x,y))$ such that

Both P and Q are not zero polynomials.

Points (x,y) in the domain \mathbb{P}^1 's that $P(x,y) = Q(x,y) = 0$ are called the base points

• Space of q -maps to \mathbb{P}^1 is $\mathbb{P}^{2(d+1)}$

Equivariant cohomology

$$H^0(\mathbb{P}^1) = \frac{\mathbb{C}[P]}{\langle p^2=0 \rangle}$$

$$H_{\mathbb{C}^*}^0(\mathbb{P}^1) = \frac{\mathbb{C}[P]}{\langle (p-\varepsilon_1)(p-\varepsilon_2) \rangle}$$

Equivariant localization

X -equivariant reduction around action on X (for simplicial has only

equivariant vars corresponding to the group action

Equivariant localization

equivariant vars corresponding to the group action

X -sm. proj, G -reductive group action on X (for simplicity has only fixed point)

$$H_G(X)_{loc} \cong H_G(X^G)_{loc}$$

↖ localized wrt equivariant vars.

$$\alpha \in H_G(X)$$

$$z(\alpha) = \sum_{p \in \text{fixed pts}} \frac{\alpha|_p}{Eu_G(T_p X)}$$

• Let's apply this to q -maps on \mathbb{P}^1

What are the fixed pts?

It has to be $(P(x,y)=0)$ or $(0:Q(x,y))$.

$$P(x,y) = \sum_{i=0}^d c_i x^i y^{d-i} \quad \xrightarrow{\quad} \quad c_i q^i x^i y^{d-i}$$

It has to be $(x^i y^{d-i} : 0)$ or $(0 : x^i y^{d-i})$

We will not allow base pts at $(0:1)$ of the domain \mathbb{P}^1

the only remaining fixed pts are $(y^d : 0)$ and $(0 : y^d)$

Let's compute $Eu_G(T_p \mathbb{P}^{2(d+1)})$

\mathbb{P}^1 $(y^d + p(x,y); q(x,y))$

$c_1 y^d + c_2 x y^{d-1} + \dots + c_{d+1} x^d$ ← deformations at the q -map.

We compute $T_p \mathbb{P}^{2(d+1)}$ as a $\mathbb{C}_q^\times \times (\mathbb{C}^\times)_{\alpha, \alpha^2}^2$ -rep.

$$T_p \mathbb{P} = \alpha_1 (1 \oplus q \oplus \dots \oplus q^d) \oplus \alpha_2 \alpha^{-1} (1 \oplus q \oplus \dots \oplus q^d)$$

$$T_p \mathbb{P}^d = \alpha_1 (1 \oplus q \oplus \dots \oplus q^d) \oplus \alpha_2 \alpha_1^{-1} (1 \oplus q \oplus \dots \oplus q^d)$$

Use the target space rescaling to set coeff in front of y^d to 1

$$Eu_G(T_p \mathbb{P}^{2(d+1)}) = \prod_{n=1}^d n \cdot z \cdot \prod_{n=0}^d (\varepsilon_2 - \varepsilon_1 + n z)$$

$$Eu_G(\oplus_i \mathcal{L}_i) = \prod_i c_i(\mathcal{L}_i)$$

let's denote $c_i(\alpha_j) = \varepsilon_i$
 $c_i(q) = z$ ← equivariant vars

$$i(z) = z^{-2d-1} \frac{\alpha_1 p_1}{d! \prod_{n=0}^d (\frac{\varepsilon_2 - \varepsilon_1}{z} + n)} + z^{-2d-1} \frac{\alpha_2 p_2}{d! \prod_{n=0}^d (\frac{\varepsilon_1 - \varepsilon_2}{z} + n)}$$

$$J^{qmap} = \sum_{d=0}^{\infty} Q^d \left[z^{-2d-1} \frac{1}{d! \prod_{n=0}^d (\frac{\varepsilon_2 - \varepsilon_1}{z} + n)} + z^{-2d-1} \frac{1}{d! \prod_{n=0}^d (\frac{\varepsilon_1 - \varepsilon_2}{z} + n)} \right]$$

[p1] [p2]

Remark K -theory class on \mathbb{P}^1 target

$$Z_D^2(\mathbb{B}) := \int_{\delta + i\mathbb{R}} d\zeta \cdot \Gamma(-\frac{\zeta}{2\pi i} + \varepsilon_1) \Gamma(-\frac{\zeta}{2\pi i} + \varepsilon_2) \cdot Q \cdot \int_{\mathbb{B}} \mathcal{F}_{\mathbb{B}}(\zeta)$$

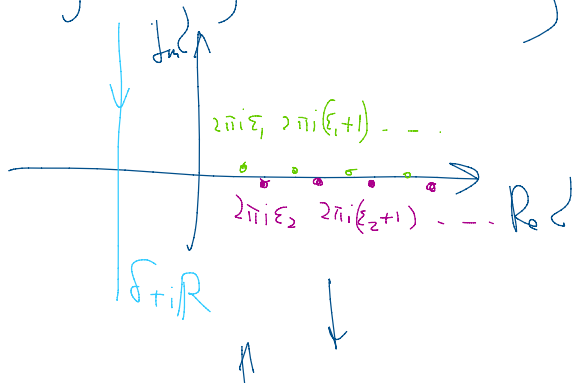
$\Gamma(x+1) = x\Gamma(x)$ $\frac{\Gamma(\varepsilon_2 - \varepsilon_1)}{z}$ explicit function

"hemisphere partition function"

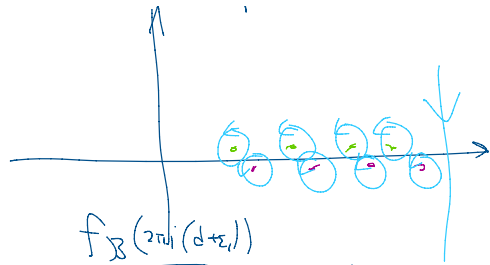
we can compute contour to $\text{Re}(\zeta) \rightarrow +\infty$

$$\text{Res}_{z=-d} \Gamma(z) = \frac{(-1)^d}{d!}$$

this integral by residues shifting the



$$Z_D^2(\beta) = \sum_{d=0}^{\infty} Q^d \left[\frac{Q^{\varepsilon_1}}{d! \Gamma(\varepsilon_1 - \varepsilon_2 + d + 1)} \frac{f_{\mathbb{P}^1}(2\pi i(d + \varepsilon_1))}{\sin \pi \varepsilon_1} + \frac{Q^{\varepsilon_2}}{\Gamma(\varepsilon_2 - \varepsilon_1 + d + 1) d!} \frac{f_{\mathbb{P}^1}(2\pi i(d + \varepsilon_2))}{\sin \pi \varepsilon_2} \right]$$



$$Z_D^2(\beta) = \left\langle \int_{z=2\pi i}^{\text{q-map}} \Gamma^{\wedge-1}, \text{ch}(\beta) \right\rangle$$

Jacobian's Γ -class
Poncare (with Todd class)

Higgs - Coulomb correspondence in our case.

Kori-Romo '13

Generalizations: target $\mathbb{P}^1 \rightsquigarrow$ toric smooth DM stack

$$[V // G] = [V^{\omega\text{-ss}} / G] = \mathcal{X}$$

$\mathbb{C}^{n+k} \xrightarrow{q} (\mathbb{C}^{\times})^k$

\mathbb{G}/\mathbb{H} quotient w.r.t. to stability parameter $\omega \in \mathfrak{g}^*$

Variable Q generalizes to e^{ω} , where $\omega \in H^2(\mathcal{X}; \mathbb{R})$

$Z_D^2(\beta)$ - analytic function in ω

Power series expansion recovers $\int^{\text{q-map}}$ - function

$\langle \int^{\text{q-map}} \Gamma^{\wedge-1}, \text{ch}(\beta) \rangle$

lower series expansion recovers

$$\langle \gamma \Big|_{z=2\pi i}^{\hat{P}^{-1}}, \text{ch}(B) \rangle$$

Expansions at diff. points recover different group theories
 q-maps to $[V //_{\omega} G]$ for different stab. ω

Example: $[\mathbb{C}^6 // \mathbb{C}^{\times}]$

$$(11111 | -5)$$

$\omega > 0$

$K_{\mathbb{P}^4}$

$\omega < 0$

$[\mathbb{C}^5 / \mathbb{Z}_5]$

version of Reid's crepant transformation

Further generalization: GLSM

"quasimap counts into critical loci of functions on $[V //_{\omega} G]$

$$\mathbb{C}^6_{x_1 \dots x_5, p}$$

$$W(\bar{x}, p) = p \cdot \sum_{i=1}^5 x_i^5$$

$$\text{Crit } W \subset \mathbb{P}^4$$

$$\text{Crit } W \subset [\mathbb{C}^5 / \mathbb{Z}_5]$$

$\cong \mathbb{S}$

$\cong \mathbb{S}$

$$\{p=0\} \cap \{\sum x_i^5 = 0\} \cong \text{quintic 3-fold in } \mathbb{P}^4$$

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