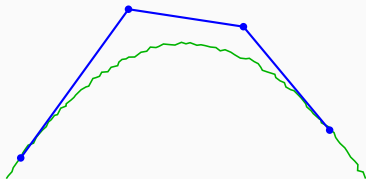


Understanding the upper tail behaviour of the KPZ equation via the tangent method

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(based on joint work with Shirshendu Ganguly)

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The Kardar-Parisi-Zhang (KPZ) equation is a non-linear stochastic PDE believed to describe planar random interface growth in a very broad class of models, and is given by

$$\partial_t H = \frac{1}{4}(\partial_x H)^2 + \frac{1}{4}\partial_x^2 H + \xi, \quad (1)$$

where ξ is space-time white noise on $\mathbb{R} \times (0, \infty)$ and $H : \mathbb{R} \times (0, \infty) \rightarrow \mathbb{R}$.

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While there are now sophisticated notions of solution available, the one that has underlied most previous studies is known as the [Cole-Hopf](#) solution.

Cole-Hopf solution to the KPZ equation

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The Cole-Hopf solution *defines* $H = \log Z$ to solve the KPZ equation, where Z solves the multiplicative stochastic heat equation

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The **narrow-wedge** solution to the KPZ equation is the solution when $Z(0, \cdot) = \delta_0$, the Dirac mass at the origin.

Scaling the narrow wedge solution

We will be interested in $H(t, x)$ for fixed $x \in \mathbb{R}$ as t varies. So it will be convenient to consider a scaled version of H which is **tight** in t .

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So we consider the **scaled** narrow-wedge solution

$$\mathfrak{h}^t(x) = \frac{H(t, t^{2/3}x) - \frac{t}{12}}{t^{1/3}};$$

this does not grow linearly with t , has unit order fluctuations, and is **tight** in $t \geq t_0 > 0$.

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It decays **parabolically**; in fact, for fixed t , $x \mapsto \mathfrak{h}^t(x) + x^2$ is a **stationary** process.



Upper tail behaviour of \mathfrak{h}^t

The **upper tail** behavior has been of significant interest in both the physics and mathematics communities.

In spite of significant recent progress relying on exact formulas available for the narrow wedge solution, the upper tail behavior is not completely understood.

Questions of interest include:

- Asymptotics of **one-** and **multi-point** tails, eg. $\mathbb{P}(\mathfrak{h}^t(0) > \theta)$ or $\mathbb{P}(\mathfrak{h}^t(-1) > \theta_-, \mathfrak{h}^t(1) > \theta_+)$.
- The **behavior of the profile** under the above events.

Existing work has been mainly focused on one-point asymptotics.

A multi-point question: Sharpness of FKG?

Because of connections to statistical mechanics models, it is known that \mathfrak{h}^t enjoys the [FKG inequality](#), so that, for all θ_- and θ_+ ,

$$\mathbb{P}(\mathfrak{h}^t(-1) > \theta_-, \mathfrak{h}^t(1) > \theta_+) \geq \mathbb{P}(\mathfrak{h}^t(-1) > \theta_-) \cdot \mathbb{P}(\mathfrak{h}^t(1) > \theta_+).$$

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But in many applications the inequality is suboptimal.

So we are led to ask: Is FKG [sharp](#) for any values of θ_- and θ_+ , and, if so, which ones?

Greater structure in the $t \rightarrow \infty$ limit

It has been long expected and recently proven (Quastel-Sarkar and Virág) that the narrow-wedge solution converges to the **parabolic Airy₂ process** \mathcal{P} as $t \rightarrow \infty$.

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This has given a sharp understanding of the upper tail asymptotics: eg., it is known (Dumaz-Virág, following work of Ramirez-Rider-Virág) that, as $\theta \rightarrow \infty$,

$$\mathbb{P}(\mathcal{P}(0) > \theta) = \exp\left(-\frac{4}{3}\theta^{3/2} + O(\log \theta)\right).$$

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The question of spatial structure has received some more attention however, with Quastel-Tsai proving a large deviation principle for a discrete prelimit (TASEP) of \mathcal{P} .

There has been some work in the case of finite t as well:

- Lamarre-Lin-Tsai investigate the upper tail large deviation limit shape for *short time* ($t \rightarrow 0$) using a Feynman-Kac representation. (In this limit the non-linearity disappears and the solution falls into the Gaussian universality class.)

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Related work in the finite t case

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The above rely on techniques such as the Feynman-Kac representation, PDE methods, and the exact formulas available for the narrow-wedge solution.

In contrast, our approach is more *geometric* and will yield near-optimal results.

Theorem (Corwin-Ghosal)

The following holds for a wide class of initial data as well as the narrow-wedge case: there exist $c_1, c_2 > 0$ and θ_0 such that, for $\theta > \theta_0$ and $t \geq 1$,

$$\exp(-c_1\theta^{3/2}) \leq \mathbb{P}(\mathfrak{h}^t(0) > \theta) \leq \exp(-c_2\theta^{3/2}).$$

The constants c_1 and c_2 are explicit but non-optimal for general initial data (predicted to be $4/3$ in the physics literature).

For narrow wedge the methods did obtain the optimal constant of $\frac{4}{3} + o(1)$, but only in certain regimes of θ .

Main results

Theorem

There exist $C < \infty$ and θ_0 such that, for all $t \geq 1$ and $\theta > \theta_0$,

$$\exp\left(-\frac{4}{3}\theta^{3/2} - C\theta^{9/8}\right) \leq \frac{1}{d\theta} \mathbb{P}(\mathfrak{h}^t(0) \in d\theta) \leq \exp\left(-\frac{4}{3}\theta^{3/2} + C\theta^{9/8}\right).$$

As an immediate consequence, the same bounds also hold for $\mathbb{P}(\mathfrak{h}^t(0) > \theta)$.

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- This gives a sharp bound with the optimal $\frac{4}{3}$ for the *density*. To our knowledge, bounds on the density were not previously available in the literature.
- The bound holds for all large values of θ with the optimal coefficient $\frac{4}{3}$, and the error is uniform in t .
- The bound also holds for \mathcal{P} (as mentioned a sharper version for the *upper tail* of \mathcal{P} was already known, but the density bound is new).
- We will give similar tail bounds for general initial data shortly.

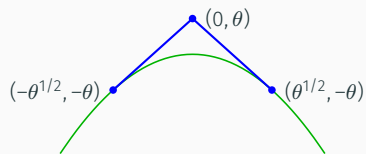
Main results: the one-point limit shape

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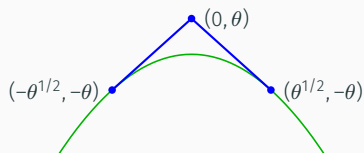


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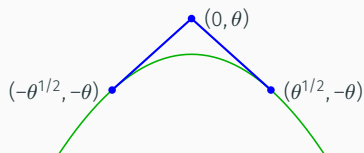
There exist θ_0 and $c > 0$ such that, for all $t \geq 1$, $\theta > \theta_0$, and $M > 0$,

$$\mathbb{P} \left(\sup_{x \in [-\theta^{1/2}, \theta^{1/2}]} |\mathfrak{h}^t(x) - \text{Triangle}_\theta(x)| > M\theta^{1/4} \mid \mathfrak{h}^t(0) > \theta \right) \leq \exp(-cM^2).$$

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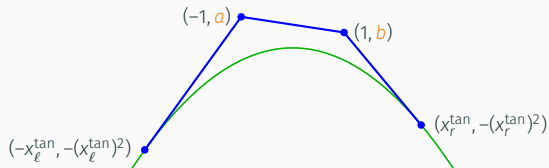
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- The bound also holds with the conditioning $\mathfrak{h}^t(0) = \theta$, and for \mathcal{P} .
- $\theta^{1/4}$ is the Brownian fluctuation scale on an interval of size $\theta^{1/2}$ and is optimal.

Main results: two-point limit shape

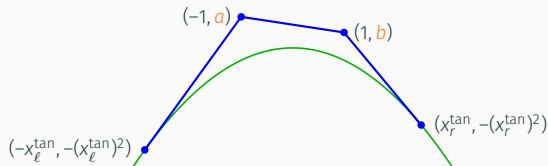
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The values of x_ℓ^{tan} and x_r^{tan} are such that the tangency conditions are met.

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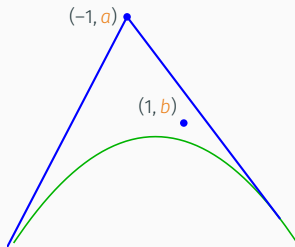
Assuming some *non-degeneracy* conditions on a and b , there exists $c > 0$ such that, for all $t \geq 1$, $M > 0$, and large enough a, b ,

$$\mathbb{P} \left(\sup_{x \in [-x_\ell^{\text{tan}}, x_r^{\text{tan}}]} |\mathfrak{h}^t(x) - \text{Quad}_{a,b}(x)| > M(a^{1/4} + b^{1/4}) \mid \mathfrak{h}^t(-1) > a, \mathfrak{h}^t(1) > b \right) \leq \exp(-cM^2).$$

- The bound again also holds for \mathcal{P} .

Main results: two-point limit shape

The **non-degeneracy** conditions are to ensure that both $(-1, a)$ and $(1, b)$ are extreme points of the convex hull, unlike below.



Theorem

For $t \geq 1$ and if a, b are large enough and satisfy the non-degeneracy condition, then

$$\begin{aligned} & \mathbb{P}(\mathfrak{h}^t(-1) \geq a, \mathfrak{h}^t(1) \geq b) \\ &= \exp\left(-\frac{1}{24} \left[16 \left((1+a)^{3/2} + (1+b)^{3/2}\right) + 3(a-b)^2 + 24(a+b) + 32\right] + \text{error}\right). \end{aligned}$$

The error term has explicit upper and lower bounds, uniformly in t . The asymptotic also holds for \mathcal{P} .

Main results: Two point asymptotics

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- This is the first sharp asymptotic for the two-point distribution we know of, and is also new for \mathcal{P} .
- The non-degeneracy condition on a, b implies that $(b-a)^2 \ll a^{3/2}, b^{3/2}$. Without it, the one-point asymptotic dominates.
- A similar bound also holds at $\pm K$ in place of ± 1 .

Corollary: Geometric condition for sharpness of FKG

Recall that the FKG inequality implies

$$\mathbb{P}\left(\mathfrak{h}^t(-K^{1/2}) > a, \mathfrak{h}^t(K^{1/2}) > b\right) \geq \mathbb{P}\left(\mathfrak{h}^t(-K^{1/2}) > a\right) \cdot \mathbb{P}\left(\mathfrak{h}^t(K^{1/2}) > b\right),$$

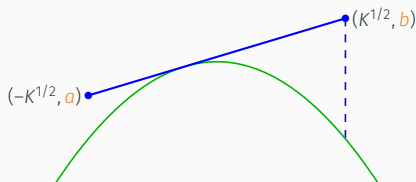
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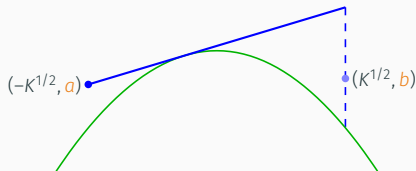


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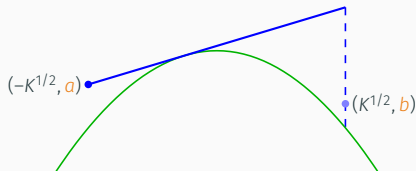


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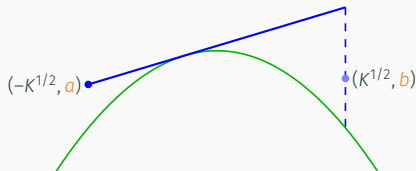


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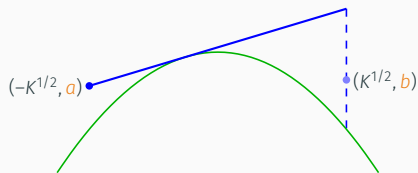
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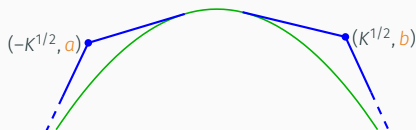
Corollary

Let K be fixed. If the line joining $(-K^{1/2}, a)$ and $(K^{1/2}, b)$ is tangent to or intersects $-x^2$ inside $[-K^{1/2}, K^{1/2}]$, then

$$\begin{aligned}\mathbb{P}\left(\mathfrak{h}^t(-K^{1/2}) > a, \mathfrak{h}^t(K^{1/2}) > b\right) &= \exp\left(-\frac{4}{3}[(K+a)^{3/2} + (K+b)^{3/2}] + \text{error}\right) \\ &\approx \mathbb{P}(\mathfrak{h}^t(-K^{1/2}) > a) \cdot \mathbb{P}(\mathfrak{h}^t(K^{1/2}) > b).\end{aligned}$$

- The second line is via the one-point asymptotics and since $\mathfrak{h}^t(\pm K^{1/2}) + K \stackrel{d}{=} \mathfrak{h}^t(0)$ by stationarity of $\mathfrak{h}^t(x) + x^2$.

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- In essence, the parabola acts as a **barrier** to the interaction of the events $\{\mathfrak{h}^t(-K^{1/2}) > a\}$ and $\{\mathfrak{h}^t(K^{1/2}) > b\}$.

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- **Not $-\infty$ everywhere:** There is a positive measure set where $h^0 \neq -\infty$.

(The hypotheses essentially ensure h^t exists for all $t \geq 1$ and is non-trivial.)

Recall that the solution h_{gen}^t is given by $\log Z^t$, where $Z^t(x)$ solves the stochastic heat equation with $Z^0(x) = \exp(h^0(x))$.

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Theorem

There exist C, θ_0 (depending on \mathfrak{h}^0) such that, for $t \geq 1$, and $\theta > \theta_0$,

$$\exp\left(-\frac{4}{3}\theta^{3/2} - C\theta^{9/8}\right) \leq \mathbb{P}\left(\mathfrak{h}_{\text{gen}}^t(0) \geq \theta\right) \leq \exp\left(-\frac{4}{3}\theta^{3/2} + C\theta^{9/8}\right).$$

The constants C and θ_0 can be made uniform over a class of initial data by quantifying the hypotheses on \mathfrak{h}^0 .

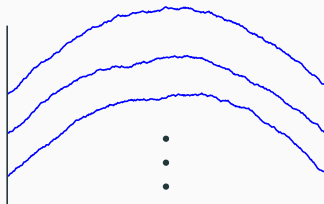
Summary of results

- The limit shapes for large one- and two-point values are given in terms of tangent lines to the parabola through the high points, and the fluctuations around the shape are Brownian.
- The information about the limit shapes can be combined with Brownian estimates to give sharp asymptotics for the one- and two-point probabilities.
- These asymptotics give a geometric understanding of asymptotic independence of upper-tail events, i.e., the sharpness of FKG.
- This can be extended to give sharp one-point asymptotics for general initial data.

The Brownian Gibbs property

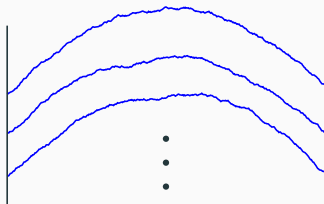
The resampling property

Both \mathfrak{h}^t and \mathcal{P} can be embedded as the top, or lowest-indexed, curve in a \mathbb{N} -indexed ensemble of random continuous curves.



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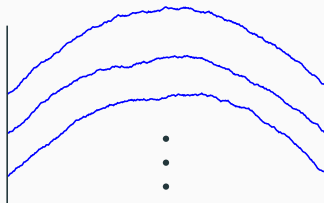


Both ensembles enjoy a resampling property known as the [Brownian Gibbs](#) property.

We focus on the simpler property for the $t = \infty$ ensemble including \mathcal{P} .

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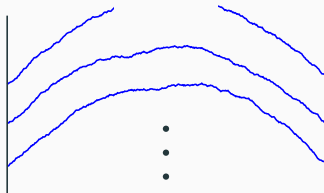
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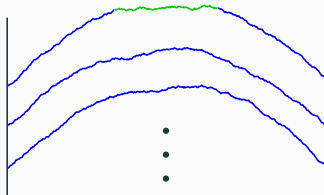
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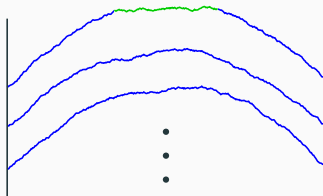
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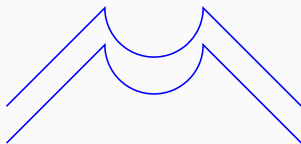
A useful heuristic to keep in mind:

\mathfrak{h}^t and \mathcal{P} are like Brownian bridges conditioned to stay above a parabola $-x^2$ with which they share endpoints.

A convex consequence

A key idea is that the resampling in terms of Brownian bridges implies that the limit shapes should be **convex**.

Indeed, suppose the limit shape of the top curve is not convex in some neighbourhood. This pushes the **second curve** down on the interval.

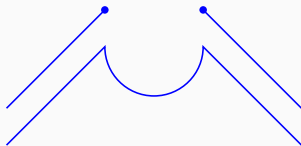


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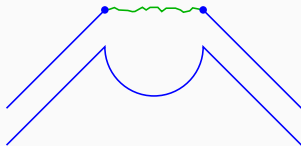
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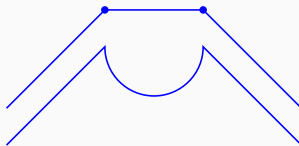
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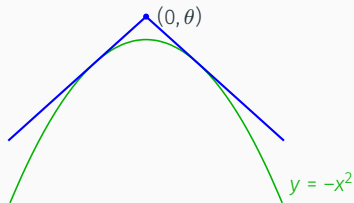


In the six-vertex model, a similar idea was termed the **tangent method** by Colomo-Sportiello and was rigorously implemented by Aggarwal.

The six-vertex model is an example of a “zero temperature” model (similar to \mathcal{P}), while the KPZ equation at finite t is a “positive temperature” model.

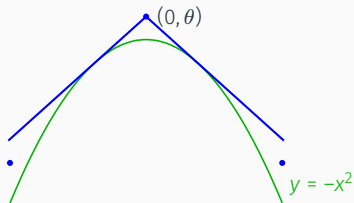
Proof ideas

Getting the one-point limit shape (upper bound): pinning



To use the linear trajectories of Brownian bridges, we first need to “pin” $\mathfrak{h}^t(x)$ to the **tangent** line at some x ; then we will know that **Triangle $_{\theta}$** is followed. Take $x = \theta^{1/2}z$.

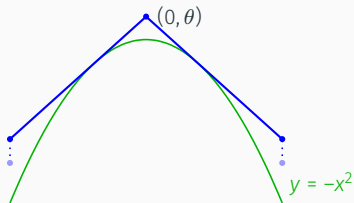
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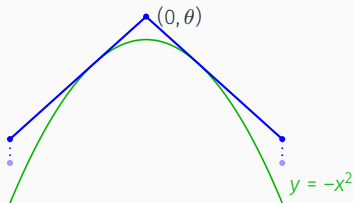
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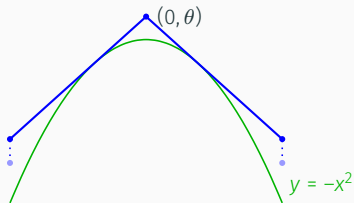
In fact to upper bound the profile shape, it's enough if \mathfrak{h}^t is **below** the tangent at some large x : we can raise the points to the tangent, and this can only raise the profile.

We need \mathfrak{h}^t to be below the tangent with high probability **conditional on $\mathfrak{h}^t(0) > \theta$** . But by stationarity + parabolic curvature,

$$\mathbb{P}\left(\mathfrak{h}^t(\theta^{1/2}z) > \text{Tangent}(\theta^{1/2}z) \mid \mathfrak{h}^t(0) > \theta\right) \leq \frac{\mathbb{P}(\mathfrak{h}^t(0) > \text{Tangent}(\theta^{1/2}z) + \theta z^2)}{\mathbb{P}(\mathfrak{h}^t(0) > \theta)},$$

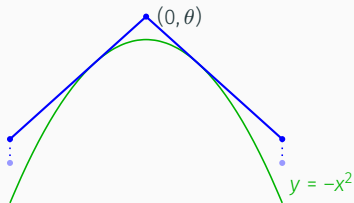
which is small for large enough z (but still $O(1)$) by the Corwin-Ghosal tail bounds.

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Once we have the pinning, the profile from $(z\theta^{1/2}, \text{Tangent}(z\theta^{1/2}))$ to $(0, \theta)$ is a Brownian bridge **conditioned** to stay above the second curve, essentially a parabola.

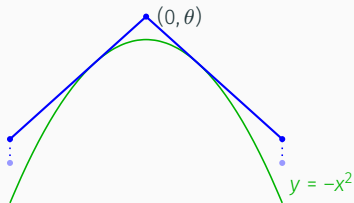
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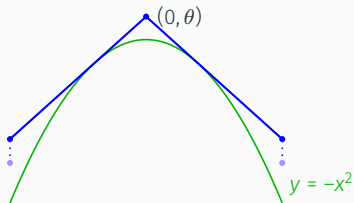


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The linear trajectory is close to $-x^2$ only at the tangency point, so Brownian bridge avoids $-x^2$ with constant probability. So the **conditioning** can essentially be ignored.

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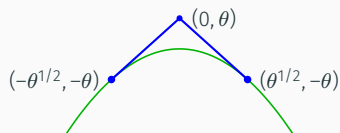
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This yields the $\exp(-cM^2)$ bound for a deviation of size M on the Brownian scale $\theta^{1/4}$.

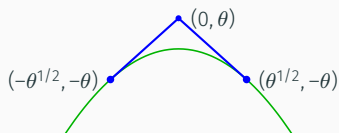
Getting the one point tail asymptotics



Since \mathfrak{h}^t looks like Triangle_θ on $[-\theta^{1/2}, \theta^{1/2}]$ when $\mathfrak{h}^t(0) > \theta$, $\mathbb{P}(\mathfrak{h}^t(0) > \theta)$ becomes a Brownian calculation: with B a (rate 2) Brownian bridge from $(-\theta^{1/2}, -\theta)$ to $(\theta^{1/2}, -\theta)$,

$$\mathbb{P}(\mathfrak{h}^t(0) > \theta) \approx \mathbb{P}\left(B(0) > \theta \mid B(x) > -x^2 \forall x \in [-\theta^{1/2}, \theta^{1/2}]\right)$$

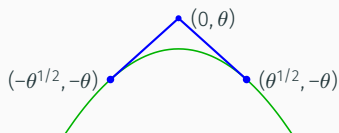
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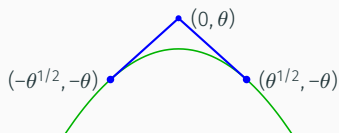
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The denominator D is $\geq \exp\left(-\frac{2}{3}\theta^{3/2}\right)$.

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$$D \leq \mathbb{P}\left(\int_{-\theta^{1/2}}^{\theta^{1/2}} B(x) + x^2 > 0\right) = \mathbb{P}\left(N\left(-\frac{4}{3}\theta^{3/2}, \frac{4}{3}\theta^{3/2}\right) > 0\right) \approx \exp\left(-\frac{2}{3}\theta^{3/2}\right),$$

since $\exp(-y^2/2y) = \exp(-y/2)$.

The argument for general initial data

A distributional convolution formula allows one to get **one-point** information for general initial data from **spatial** information for narrow-wedge:

$$h_{\text{gen}}^t(0) \stackrel{d}{=} t^{-1/3} \log \int_{\mathbb{R}} \exp \left[t^{1/3} \left\{ h_{\text{nw}}^t(x) + h^0(x) \right\} \right] dx.$$

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This can be done with further resampling arguments using the Brownian Gibbs property.

- Using geometric methods combined with the Brownian Gibbs properties, we can obtain **sharp asymptotics for one- and two-point upper tails** for narrow-wedge solutions, as well as for one-point asymptotics for general initial data.
- As a first step for this, and also for independent interest, we need to understand the shape of the profile under these asymptotic events, which we do using ideas similar to the **tangent method**.
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Thank you!