

# ODE 527

- Goals:
- 1) Finish case  $r_1, -r_2 \in \mathbb{Z}$  of the Frobenius method
  - 2) Singular points at  $\infty$ : change of variables & why hypergeometric functions show up everywhere  
 ${}_2F_1(\alpha, \beta; \gamma; x)$
  - 3) Chebyshev polynomials  $\rightsquigarrow$  Bessel functions

Recall: The Frobenius method for solving singular ODEs

- $n$ th order similarly

$$P(x)y'' + Q(x)y' + R(x)y = 0$$

$$x^2 y'' + x p(x)y' + q(x)y = 0$$

- $p(0)$  well defined
- $p(x)$  analytic
- $q(0)$  well defined &
- $q(x)$  analytic

Frobenius ansatz

indicial equation  
characteristic exponents

$$r(r-1) + r p(0) + q(0) = 0 \quad \text{has root } r_1 \text{ and } r_2$$

The ansatz

$$y_1 = x^{r_1} \sum_{n=0}^{\infty} a_n x^n$$

$$y_2 = x^{r_2} \sum_{n=0}^{\infty} a_n x^n \quad r_2 - r_1 \notin \mathbb{Z}$$

If  $r_1 - r_2$  is an integer: we may not get a second solution because we're basically making the same guess.

$r_1 < r_2$ , we can find

$$y_1 = x^{r_1} \sum_{n=0}^{\infty} a_n x^n \quad \text{as a solution}$$

if  $k = r_2 - r_1$ , then our ansatz would be

$$y_2 = x^{r_2} \sum_{n=0}^{\infty} b_n x^n = x^k x^{r_1} \sum_{n=0}^{\infty} b_n x^n = x^{r_1} \sum_{n=0}^{\infty} b_n x^{n+k}$$

We have to try something different in this case

## Frobenius ansatz if $r_2 - r_1 \in \mathbb{Z}$

If  $r_1$  &  $r_2$  are roots of indicial equation

$$r_2 < r_1 \quad r_2 - r_1 \in \mathbb{Z}$$

Ansatz

$$y_1 = x^{r_1} \sum_{n=0}^{\infty} a_n x^n$$

$$y_2 = c y_1(x) \log(x) + x^{r_2} \sum_{n=0}^{\infty} b_n x^n \quad c \in \mathbb{R}$$

Ex: Euler equation

$$x^2 y'' + \alpha x y' + \beta y = 0 \quad \alpha, \beta \in \mathbb{R}$$

$$I(r) = r(r-1) + \alpha r + \beta = 0$$

$$r^2 + (\alpha-1)r + \beta = 0$$

Ex:  $\beta = 4$   $\alpha = -3$

$$r^2 - 4r + 4 = 0$$

$$(r-2)^2 \quad r_1 = 2, r_2 = 2$$

$r_1 - r_2$  is an integer.

We learned that

$y_1 = x^{r_1}$  is a solution

(last line:  $r_1 - r_2 \in \mathbb{Z}$ )  
 $(c_1 x^{r_1} + c_2 x^{r_2})$

Our ansatz is

$$y_2(x) = c y_1(x) \log(x) + x^{r_2} \sum_{n=0}^{\infty} b_n x^n$$

in this case:  $b_n = 0$  for  $n = 0, 1, \dots$

Solution is

$$y_2(x) = y_1(x) \log(x) = x^{r_1} \log(x) = x^2 \log(x)$$

$$= \log(x) e^{2 \log(x)} \sim t e^{2t} \quad \text{if } t = \log(x)$$

In general: We are solving a Euler equation with varying  $\alpha, \beta$   
so we need the  $x^{r_2} \sum_{n=0}^{\infty} b_n x^n$  as well.

Last time: The hypergeometric function

$${}_2F_1(\alpha, \beta | \gamma; x) = \sum_{n=0}^{\infty} \frac{(\alpha)_n (\beta)_n}{n! (\gamma)_n} x^n$$

$(a)_n = a(a+1)\dots(a+n-1)$   
is rising factorial,  
Pochhammer symbol

- solution around  $x=0$  to

$$x(1-x)y'' + (\gamma - (1+\alpha+\beta)x)y' - \alpha\beta y = 0$$

$r_1=0$   $r_2=1-\gamma$  if  $r_2-r_1 \notin \mathbb{Z}$  we have other solution

$$x^{1-\gamma} {}_2F_1(\alpha, \beta | \gamma'; x)$$

if  $r_2-r_1 \in \mathbb{Z}$   
 $(\log(x)) {}_2F_1(\alpha, \beta | \gamma; x) +$  more complicated series.

Claim: The hypergeometric ODE has regular singular points at  $0, 1, \infty$

0: What we did before

$$1: x(1-x)y'' + (\gamma - (1+\alpha+\beta)x)y' - \alpha\beta y = 0$$

$$\lim_{x \rightarrow 1} (x-1) \frac{Q(x)}{P(x)} = \lim_{x \rightarrow 1} \frac{(x-1) \cdot (\gamma - (1+\alpha+\beta)x)}{x(1-x)} = \gamma - 1 - \alpha - \beta \quad \checkmark$$

$$\lim_{x \rightarrow 1} (x-1)^2 \frac{R(x)}{P(x)} = \frac{-(x-1)^2 \alpha \beta}{x(1-x)} = 0 \quad \checkmark$$

Q: How do we check at  $\infty$ ?

A: Change of variables  $\xi = \frac{1}{x}$   $\lim_{x \rightarrow \infty} \frac{1}{x} = 0$  study  $x=\infty$  by studying  $\xi=0$

First: Simpler example the ODE for  $\sinh(x)$  and  $\cosh(x)$

$y'' = y$  vs  $y'' - y = 0$  - Looks like it has no singular points (constant coefficients!) but actually at  $x=\infty$  it does

$$y''(x) - y(x) = 0 \quad \frac{d\xi}{dx} = \frac{1}{x^2} \quad \frac{d}{dx} = \frac{d\xi}{dx} \frac{d}{d\xi} = \frac{-1}{x^2} \frac{d}{d\xi}$$

$$\frac{dy}{dx} = \frac{d\xi}{dx} \frac{dy}{d\xi} = \frac{-1}{x^2} \frac{dy}{d\xi}$$

$$\begin{aligned} \frac{d^2 y}{dx^2} &= \frac{d}{dx} \left( \frac{dy}{dx} \right) = \frac{d}{dx} \left( \frac{-1}{x^2} \frac{dy}{d\xi} \right) = \frac{d}{dx} \left( -\xi^{-2} \frac{dy}{d\xi} \right) \\ &= \left( -\xi^{-2} \frac{d}{d\xi} \right) \left( -\xi^{-2} \frac{dy}{d\xi} \right) \\ &= -\xi^{-2} \left( -2\xi^{-3} \frac{dy}{d\xi} - \xi^{-2} \frac{d^2 y}{d\xi^2} \right) \\ &= 2\xi^{-5} \frac{dy}{d\xi} + \xi^{-4} \frac{d^2 y}{d\xi^2} \end{aligned}$$

$$y''(x) - y(x) = 2\xi^{-5} y'(\xi) + \xi^{-4} y''(\xi) - y(\xi) = 0$$

For this equation

$\xi = 0$  is an irregular singular point

$$\lim_{\xi \rightarrow 0} \xi \frac{Q(\xi)}{P(\xi)} = \lim_{\xi \rightarrow 0} \frac{\xi^5}{2\xi^3} = 0 \quad \checkmark$$

$$\lim_{\xi \rightarrow 0} \xi^2 \frac{R(\xi)}{P(\xi)} = \lim_{\xi \rightarrow 0} \frac{-\xi^2}{2\xi^3} \rightarrow -\infty \quad \times$$

This is an irregular singular point: reflected in the fact that

$\cosh(x)$  has really large asymptotics at  $x \rightarrow \infty$ , i.e.

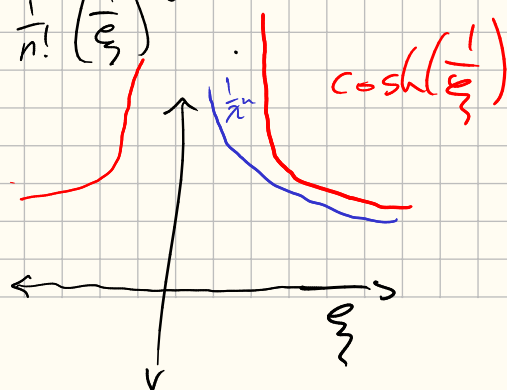
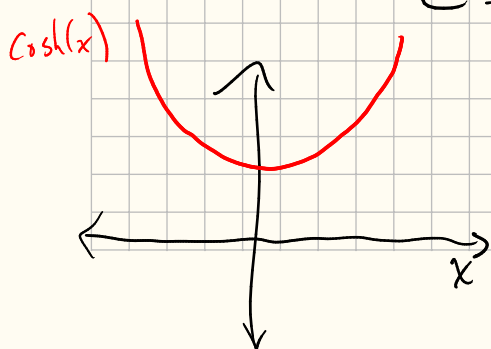
$\cosh\left(\frac{1}{\xi}\right)$  has bad asymptotics at  $\xi = 0$

$$\cosh(x) = \frac{e^x + e^{-x}}{2}$$

$$\cosh\left(\frac{1}{\xi}\right) = \frac{e^{\frac{1}{\xi}} + e^{-\frac{1}{\xi}}}{2} \rightarrow \infty \text{ at } \xi = 0 \text{ faster than any polynomial}$$

$$\lim_{\xi \rightarrow 0} \xi^N \cosh\left(\frac{1}{\xi}\right) = \infty \text{ for all } N$$

because  $e^{\frac{1}{\xi}} = \sum_{n=0}^{\infty} \frac{1}{n!} \left(\frac{1}{\xi}\right)^n$



We could never find a solution

$$\cosh\left(\frac{1}{\xi}\right) = \xi^r \sum_{n=0}^{\infty} a_n \xi^n \quad \text{using the Frobenius method.}$$

This behavior doesn't happen for  ${}_2F_1(\dots)$  after change of variables

Check:

$x = \infty$  is a regular singular point of

$$x(1-x)y'' + (\gamma - (1+\alpha+\beta)x)y' - \alpha\beta y = 0$$

$$\xi = \frac{1}{x} \quad x = \frac{1}{\xi}$$

$$\frac{dy}{dx} = -\xi^2 \frac{dy}{d\xi}$$

$$\frac{d^2y}{dx^2} = \xi^4 \frac{d^2y}{d\xi^2} + 2\xi^3 \frac{dy}{d\xi}$$

$$\frac{1}{\xi} \left(1 - \frac{1}{\xi}\right) \left(\xi^3 y''(\xi) + 2\xi^2 y'(\xi)\right) - \left(\gamma - (1+\alpha+\beta)\frac{1}{\xi}\right) \xi^2 y'(\xi) - \alpha\beta y = 0$$

$$\left(1 - \frac{1}{\xi}\right) y'' - \left[\left(1 - \frac{1}{\xi}\right) 2\xi - \left(\gamma - (1+\alpha+\beta)\frac{1}{\xi}\right) \xi^2\right] y'(\xi) - \alpha\beta y = 0$$

$$\xi=0 \quad \lim_{\xi \rightarrow 0} \frac{\xi Q(\xi)}{P(\xi)} = \frac{\xi \xi^2 (\text{finite})}{(\text{finite})} \rightarrow \text{finite} \quad \checkmark$$

$$\lim_{\xi \rightarrow 0} \frac{\xi^2 R(\xi)}{P(\xi)} = \frac{\xi^2 \alpha\beta}{1 - \frac{1}{\xi}} \rightarrow 0 \quad \checkmark$$

Conclusion  $0, 1, \text{ and } \infty$  are the only singular points & they are regular

Fact: Any 2nd order linear ODE with 3 regular singular pts (includes  $\infty$ ) is equivalent to the hypergeometric ODE after a change of variables given by a Möbius transformation

$$x \mapsto \frac{ax+b}{cx+d} \quad a, b, c, d \text{ constants.}$$

Remark: Given  $x_1, x_2, x_3$  the transformation

$$1) \quad x \mapsto \frac{(x-x_1)(x_2-x_3)}{(x-x_3)(x_2-x_1)}$$

sends  $x_1, x_2, x_3 \mapsto 0, 1, \infty$

2) Essentially:  $\infty$  being a regular singular point  $\Rightarrow$   
 $P(x)y'' + Q(x)y' + R(x)y = 0$  has low-degree polynomials for  $P(x), Q(x), R(x)$

Eg: Chebyshev (and other) orthogonal polynomials

see: demo

Chebyshev eq:  $(1-x^2)y'' - xy' + n^2y = 0$

Regular singular points at  $1, -1, \infty$   $\leftarrow$  (check)

Write the Chebyshev equation as hypergeometric eq. after Möbius transformation

$$1 \mapsto 0; \quad -1 \mapsto 1; \quad \infty \mapsto \infty$$

Möbius transformation  $x \mapsto \frac{1}{2}(1-x) = \frac{-\frac{1}{2}x + \frac{1}{2}}{0x + 1}$

$$\frac{d}{dx} = \frac{dx'}{dx} \frac{d}{dx'} = -2 \frac{d}{dx'}$$

$$x(1-x)y'' - (x-\frac{1}{2})y' + n^2y = 0$$

This is the Hypergeometric eq with  $\gamma = \frac{1}{2}, \alpha = n, \beta = -n$

Therefore  ${}_2F_1(n, -n | \frac{1}{2}; \frac{1}{2}(1-x))$  is a solution to the

Chebyshev equation (actually the polynomial solutions)

$$T_n(x) = {}_2F_1(n, -n | \frac{1}{2}; \frac{1}{2}(1-x))$$

Recall:  $T_n(\cos(\theta)) = \cos(n\theta)$

$$\Rightarrow {}_2F_1(n, -n | \frac{1}{2}; \frac{1}{2}(1-\cos(\theta))) = \cos(n\theta)$$

$$\sum_{k=0}^n \frac{\binom{n}{k} (-k)_n}{n! \left(\frac{1}{2}\right)_n} \left(\frac{1}{2} - \frac{1}{2} \cos(\theta)\right)^n = \cos(n\theta)$$

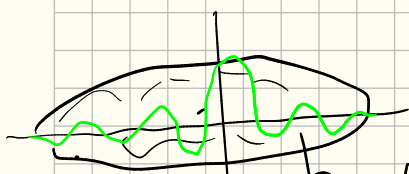
One other application: (one approach to HW#6, Problem 4)

Application to Bessel functions: solutions to

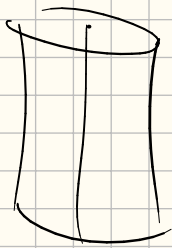
$$x^2 y'' + x y' + (x^2 - \nu^2) y = 0$$

Bessel equation of order  $\nu$ .

This arises in wave dynamics in cylindrical/spherical coordinates



Bessel functions



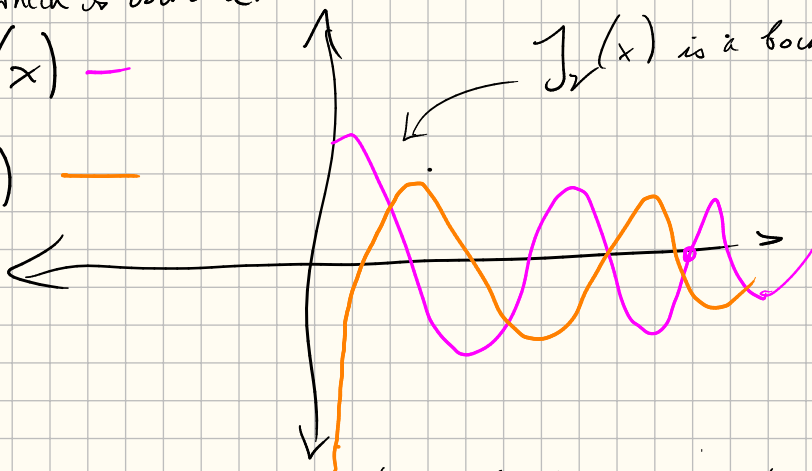
There is one solution around regular singular pt  $x=0$  which is bounded:

Bessel functions of the first kind

$$J_\nu(x)$$

$$Y_\nu(x)$$

second kind



$J_\nu(x)$  is a bounded solution

The Bessel equation has characteristic exponents  $\nu, -\nu$  (assume  $\nu \geq 0$ )

Always one bounded solution

$$J_\nu(x) = x^\nu \cdot \sum a_n x^n = \text{Bessel function of the first kind}$$

Two cases:  $\nu_2 - \nu_1 \in \mathbb{Z}$

$J_{-\nu}(x)$  is the other solution, but we take a different linear combination

$$Y_\nu(x) = \frac{J_\nu(x) \cos(\nu\pi) - J_{-\nu}(x)}{\sin(\nu\pi)} = \text{Bessel function of the second kind}$$



Other case:  $r_1 - r_2 \in \mathbb{Z}$  (i.e.  $\nu \in \frac{1}{2}\mathbb{Z}$ )

$$\begin{aligned} Y_n(x) &= y_2^{(1)} \text{ from Frobenius method} \\ &= c(\log(x)) J_n(x) + x \sum_{k=0}^{\infty} \dots \\ &= \lim_{\nu \rightarrow n} \frac{J_\nu(x) \cos(\nu\pi) - J_{-\nu}(x)}{\sin(\nu\pi)} \end{aligned}$$

Other approach:

$J_\nu(x) = \text{limit of } {}_2F_1(\alpha, \beta | \gamma; x)$  where the regular points @ 1 &  $\infty$  collide

Bessel eq. has an irregular singular pt @  $\infty$

Ex:  $\lim_{\beta \rightarrow \infty} {}_2F_1(\alpha, \beta | \gamma; \frac{x}{\beta}) = {}_1F_1(\alpha | \gamma; x)$  - irregular singular pt @  $\infty$

$$\lim_{\alpha \rightarrow \infty} {}_1F_1(\alpha | \gamma; \frac{x}{\alpha}) = {}_0F_1(1 | \gamma; x)$$

$$J_\nu(x) = c x^\nu {}_0F_1(1 | \nu + 1; -\frac{1}{4}x^2)$$

- This is a great way to write  $J_\nu(x)$  as a power series.

(see demo)